

Generating Large-scale Magnetic Fields from Cosmological Density Fluctuations

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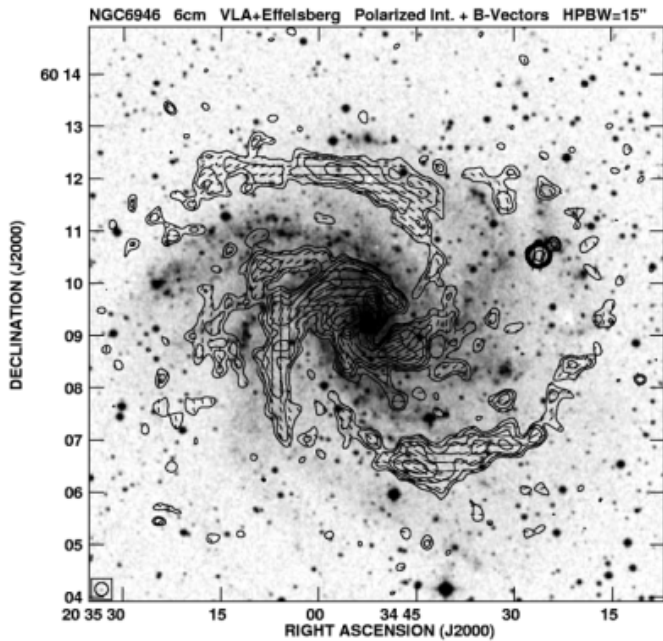
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Contents

- Introduction
 - existence of large scale magnetic fields
 - **magnetic field generation from cosmological fluctuations**
(K.I, Takahashi, Ohno, Hanayama, Sugiyama, science, (2006))
- Toward a detection of weakest cosmological magnetic fields
 - **using time-delays of emissions from γ -ray burster**
(K.I, Inoue, Takahashi, (to be submitted))
- summary

Cosmic Magnetic Fields (what is the origin?)

coherent magnetic fields at large scales

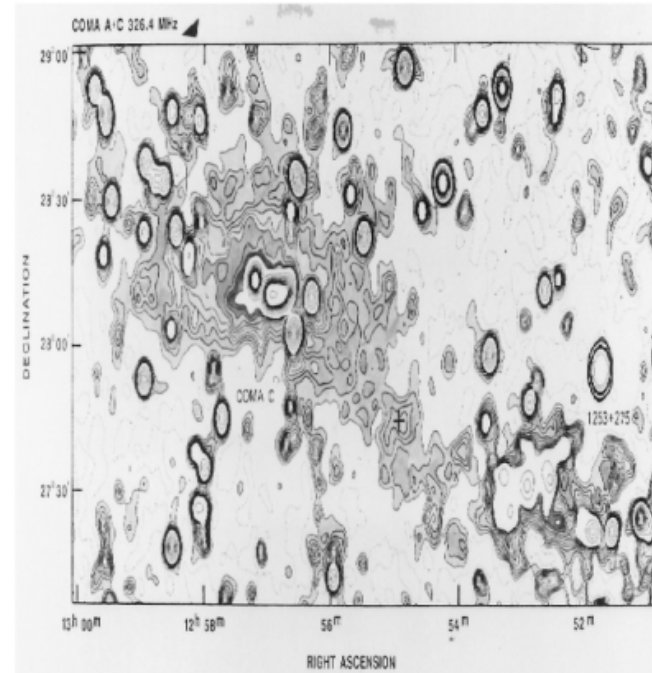


Beck and Hoernes, Nature, (1996)

NGC 6946

$B \sim \mu\text{G}$ may be associated with density spiral

structure.



Kim et al., Nature, (1989)

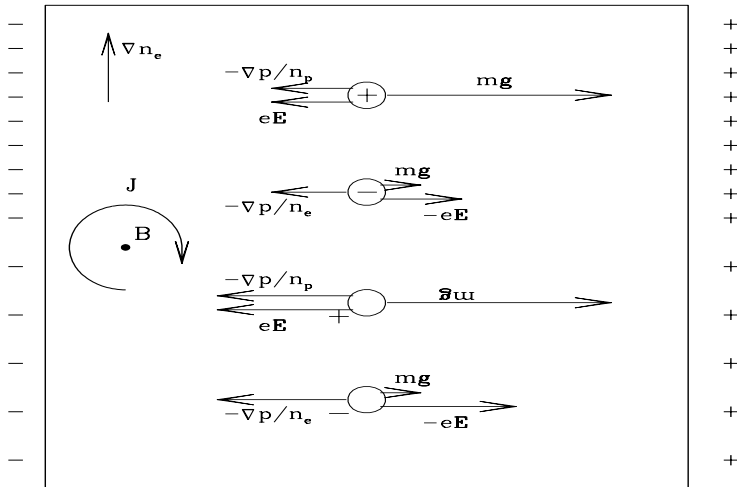
1.5Mpc "bridge" between clusters from Coma

to Abell 1367

$0.2 \sim 0.6\mu\text{G}$ (derived by minimum energy)

What is the Origin(s)? – astrophysics

- Astrophysical activities during structure formation until today
 - Biermann Battery (Biermann 1950) + dynamo amplification
 - primordial SNs: $B \sim 10^{-16} \text{G}$ (Hanayama et al, 2005)
 - protogalaxy formation: $B \sim 10^{-21} \text{G}$ (Kulsrud et al., 1997)
 - Reionization: $B \sim 10^{-18} \text{G}$ (Gnedin et al, 2000)
 - Weibel instability in structure formation shocks: $B \sim 10^{-8} \text{G}$ (Fujita et al., 2005)



- Well known physics
- strong enough
- can not predict their spectrum
 - coherence length?
- many uncertainties about predictions

What is the Origin(s)? – Cosmological

- magnetic field generation during inflation
 - need to break conformal invariance of EM fields

- $\left(\frac{\partial^2}{\partial\eta^2} - \nabla^2\right)(a^2\vec{B}) = 0 \rightarrow \vec{B} \propto a^{-2}$

- couplings to gravity: $R_{\mu\nu}A^\mu A^\nu, RF^{\mu\nu}F_{\mu\nu}$

- $B \sim 10^{-9}\text{G}$ (Turner, 1988)

- coupling to inflaton: $e^\phi F^{\mu\nu}F_{\mu\nu}$

- $10^{-65} \lesssim B \lesssim 10^{-10}\text{G}$ (Ratra, 1992)

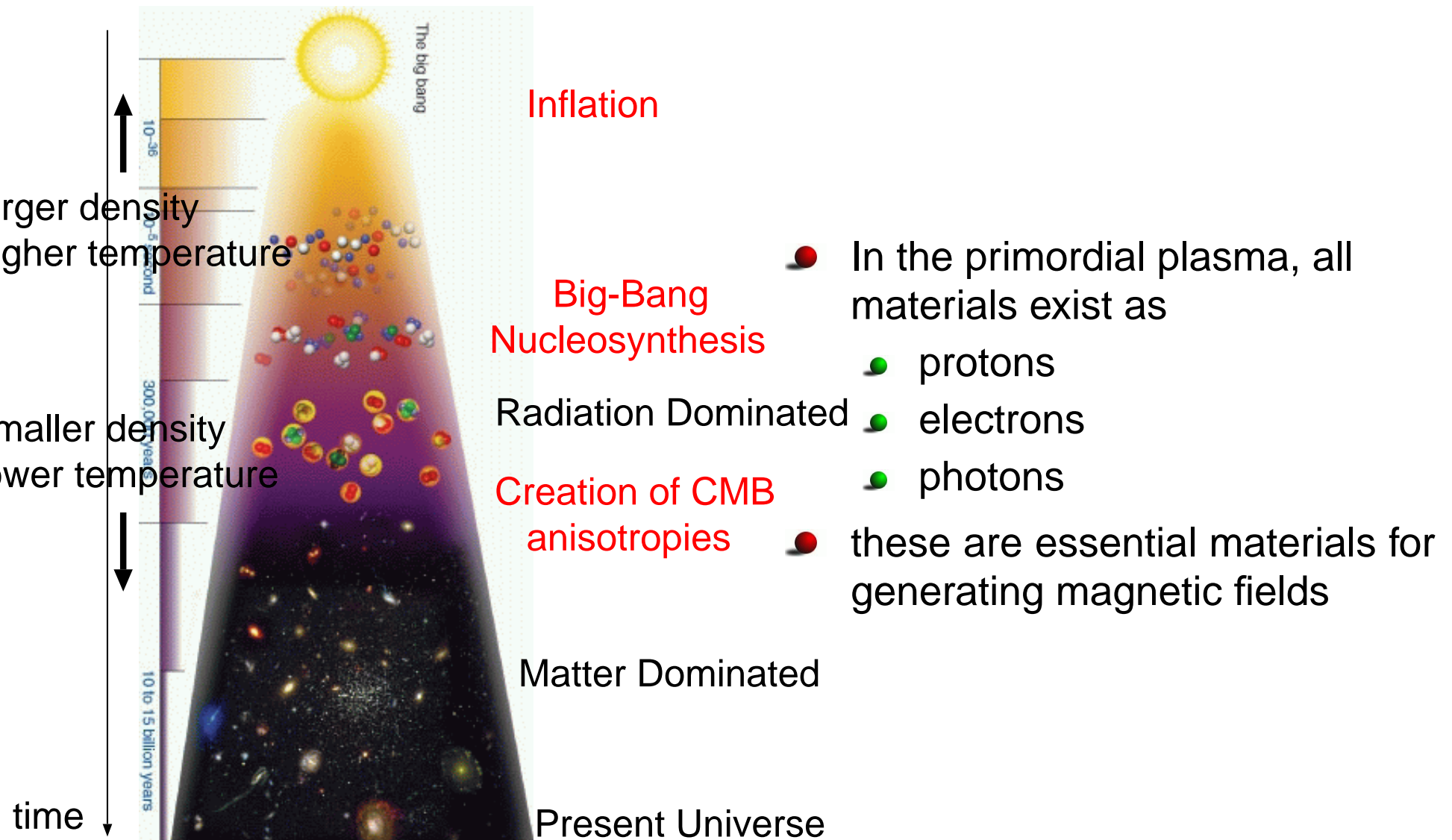
- strength and spectrum of mag. fields are predictable
- results are strongly model dependent
(which is correct?)

Our proposal – Cosmo-Physical approach

- magnetic field generation *before recombination*.
 - photons push electrons preferentially to protons **by Compton scattering** ($m_e \ll m_p$)
 - difference in fluid velocity between photons and baryons ($v_\gamma - v_b$) induces that between protons and electrons. (Harrison (1970), Gopal&Sethi(2004))
 - anisotropic stress of photons (Π_γ) pushes electrons in a different way from protons **(new!!)**
→ electric current → magnetic fields
- ($v_\gamma - v_b$), and Π_γ originate from density fluctuations.
→ **no assumption needed for existence.**

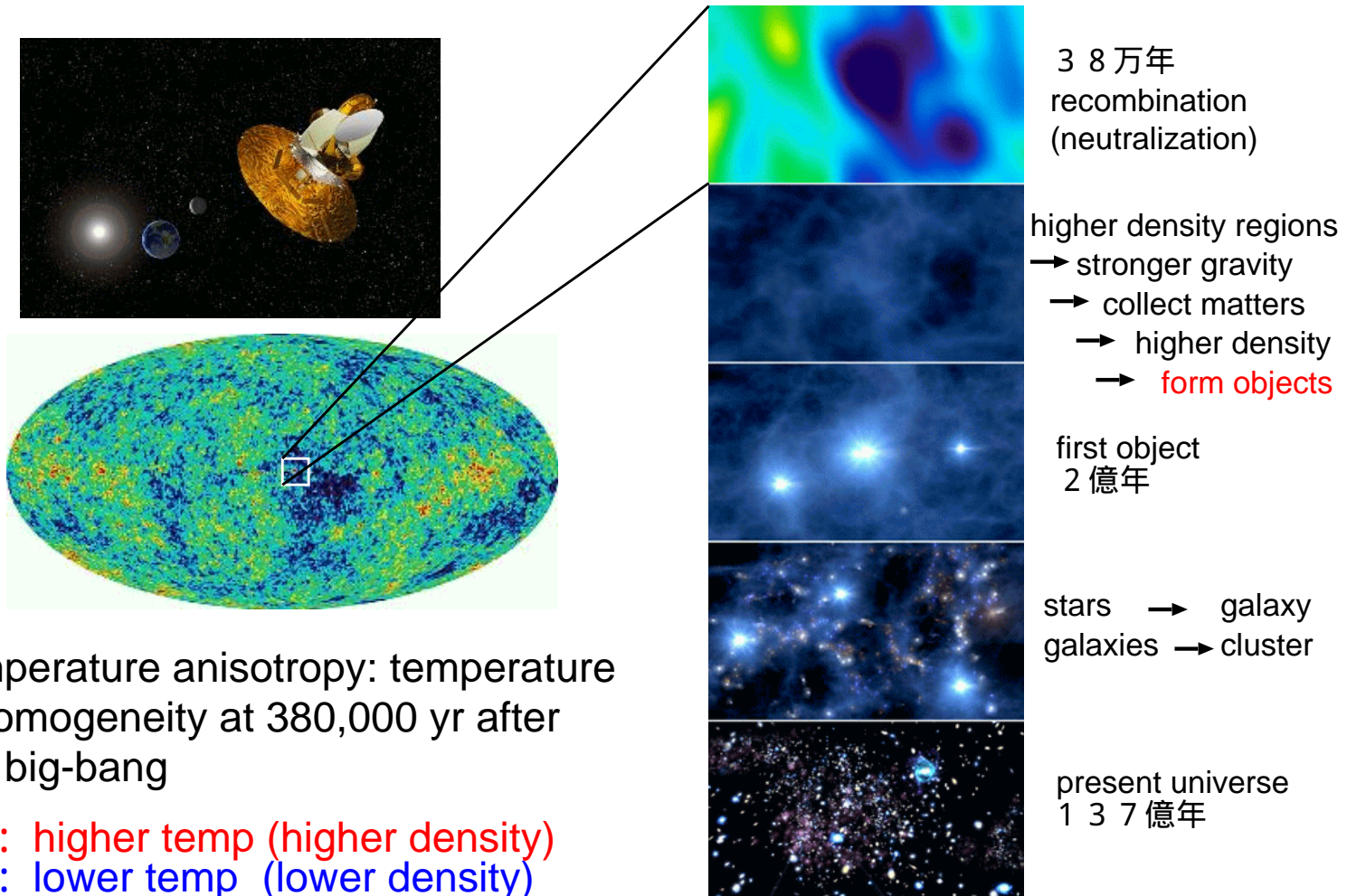
- Based on well-established cosmological perturbation theory (and observations)
- large coherence length
- field strength is small ($10^{-30} \text{G} \sim 10^{-20} \text{G}$)

History of the universe



Density Fluctuation: 'Seed' for every structures

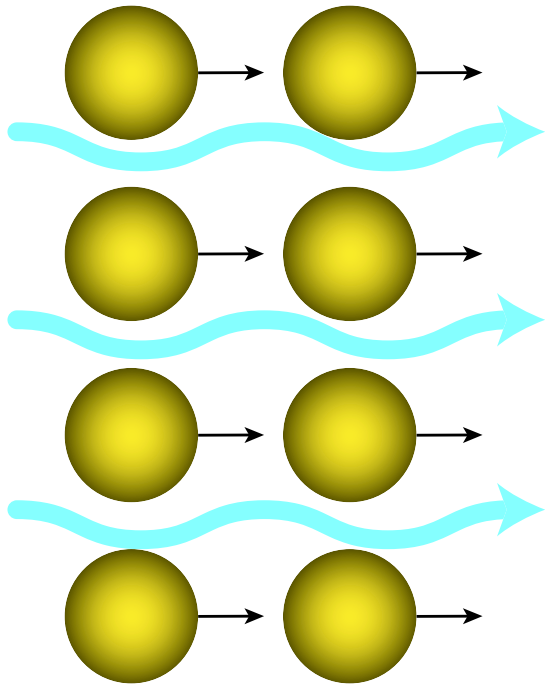
Although the universe is almost isotropic and homogeneous, we have small inhomogeneities as perturbations on it.



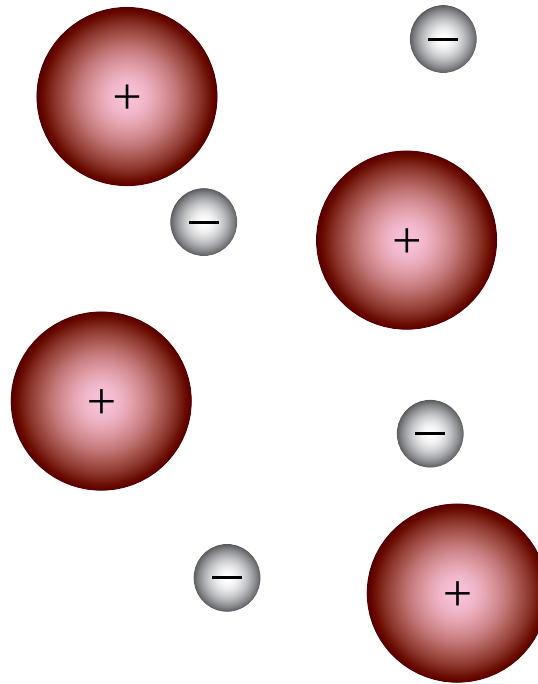
magnetic fields generation from density fluctuations

Density fluctuation of photons

-> 'photon wind' blows from high density regions



photon 'wind'

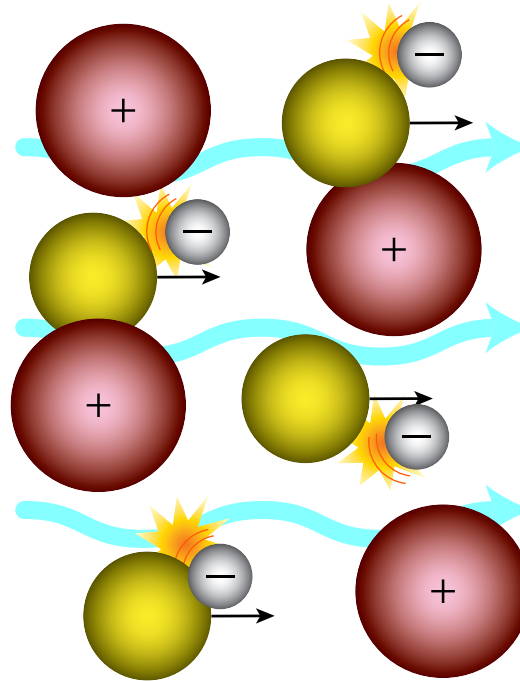
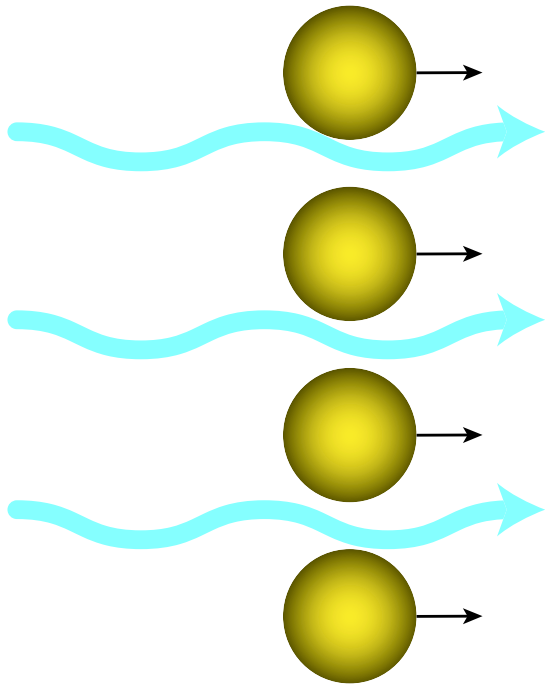


primordial plasma

magnetic fields generation from density fluctuations

Density fluctuation of photons

-> 'photon wind' blows from high density regions

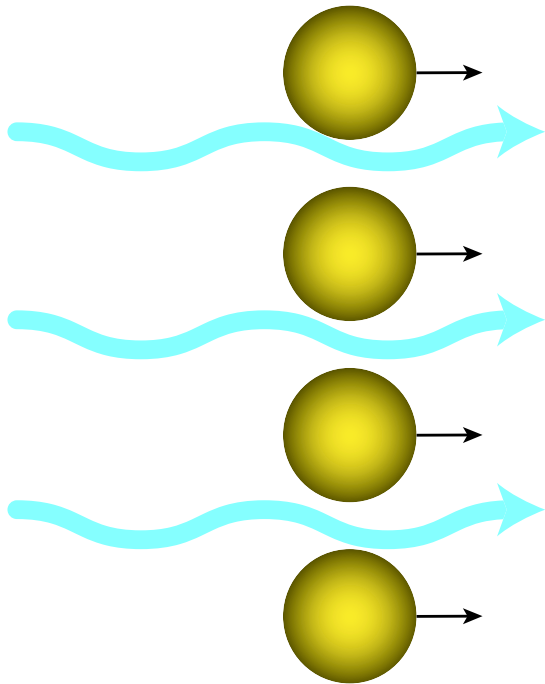


photons push only electrons

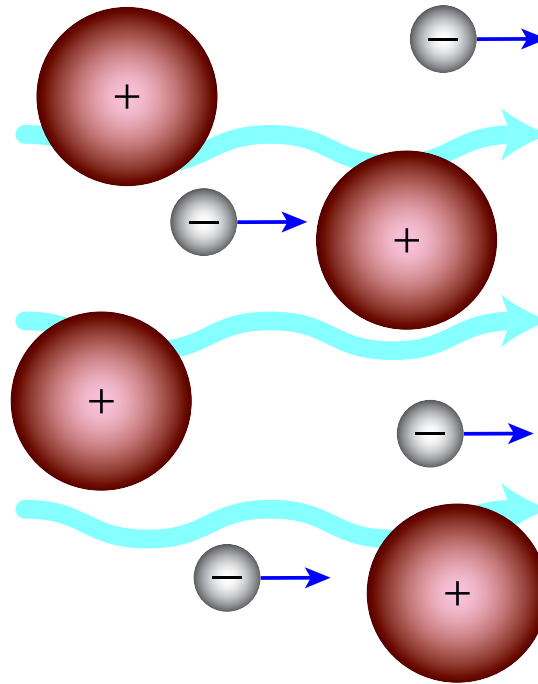
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Density fluctuation of photons

-> 'photon wind' blows from high density regions



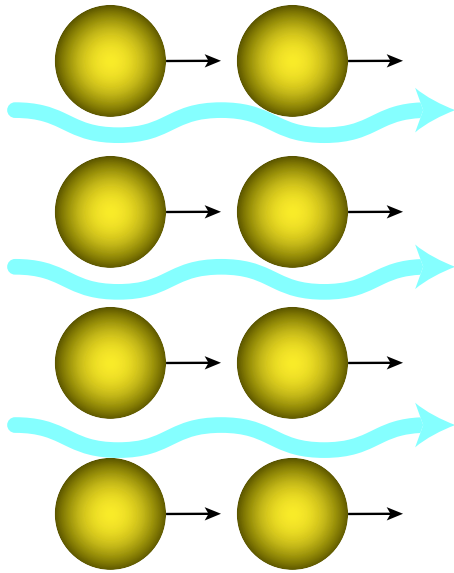
photon 'wind'



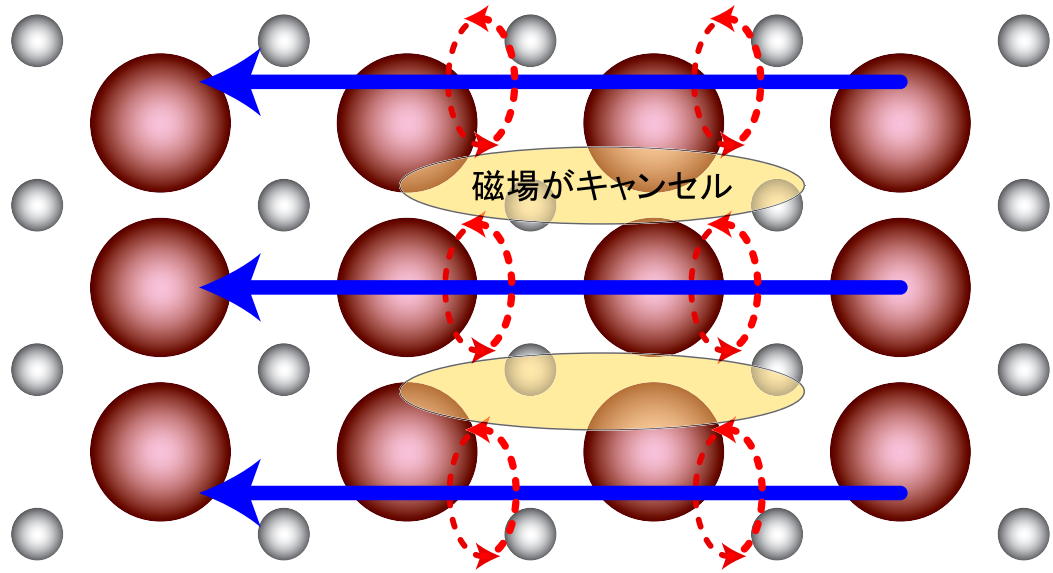
electric current -> mag fields

An Essence to Generate Mag Fields

If you work in the standard 'linear' perturbation theory...



photon 「wind」

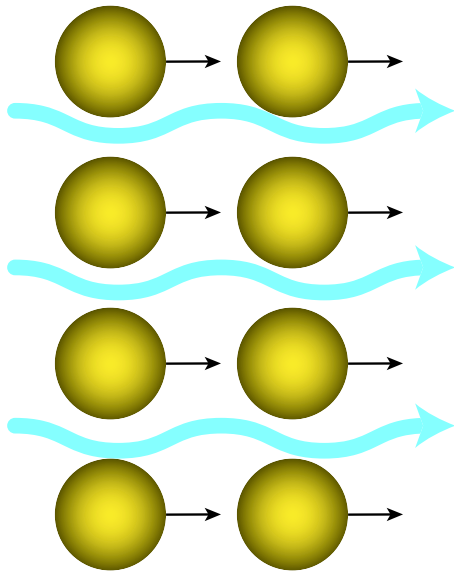


homogenous matter distribution

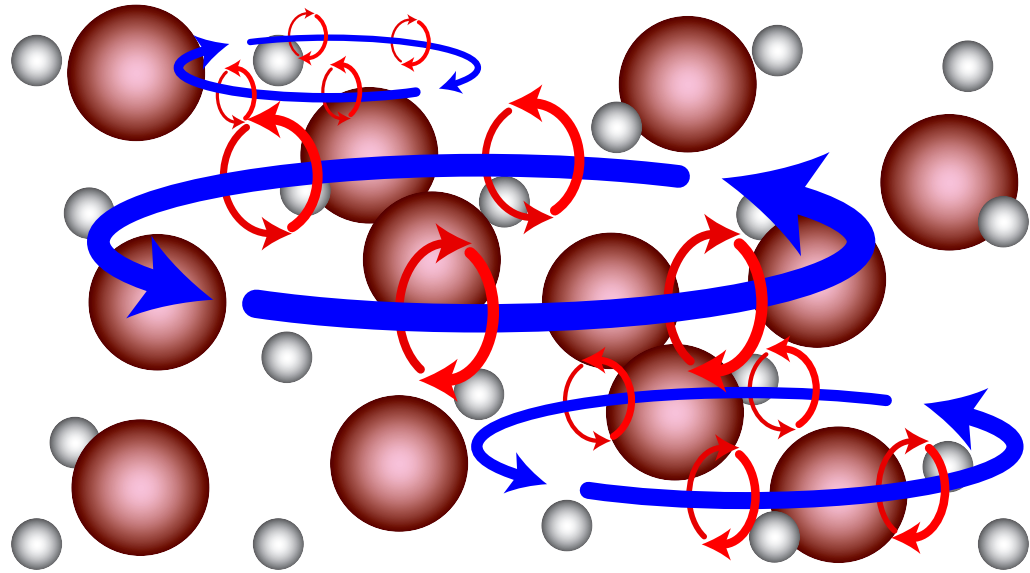
● No magnetic fields can be generated

An Essence to Generate Mag Fields

If you consider second order couplings into account...



photon 「wind」



inhomogeneous matter distribution

- magnetic field should be generated
 - in particular, anisotropic distribution of photons is important
- coupling between photon wind and inhomogeneous matter distribution (second order effect)

Evolution equation for Magnetic Fields

Vorticity difference: intrinsically second order

$$\begin{aligned}
 \dot{B}^i &= -\epsilon^{ijk} C_{j,k} \\
 &= \frac{4\sigma_T \rho_\gamma^{(0)}}{3e} \epsilon^{ijk} \left[\delta_{\gamma,k}^{(1)} \left(u_{ej}^{(1)} - u_{\gamma j}^{(1)} \right) + \left(u_{ej,k}^{(2)} - u_{\gamma j,k}^{(2)} \right) \right. \\
 &\quad \left. + \frac{1}{8} \left(u_{el,k}^{(1)} \Pi_{\gamma j}^{(1)l} + u_{el}^{(1)} \Pi_{\gamma j,k}^{(1)l} \right) \right],
 \end{aligned}$$

 : photon push

 : anisotropic pressure of photons

We can evaluate (1st order) \times (1st order) terms by non-linear convolution of perturbation spectra.

Evolution equation for Magnetic Fields

Vorticity difference: intrinsically second order

$$\begin{aligned} \dot{B}^i &= -\epsilon^{ijk} C_{j,k} \\ &= \frac{4\sigma_T \rho^{(0)} \gamma}{3e} \epsilon^{ijk} \left[\delta_{\gamma,k}^{(1)} \left(u_{ej}^{(1)} - u_{\gamma j}^{(1)} \right) + \left(u_{ej,k}^{(2)} - u_{\gamma j,k}^{(2)} \right) \right. \\ &\quad \left. + \frac{1}{8} \left(u_{el,k}^{(1)} \Pi_{\gamma j}^{(1)l} + u_{el}^{(1)} \Pi_{\gamma j,k}^{(1)l} \right) \right], \end{aligned}$$

ALL TERMS are at SECOND ORDER!!

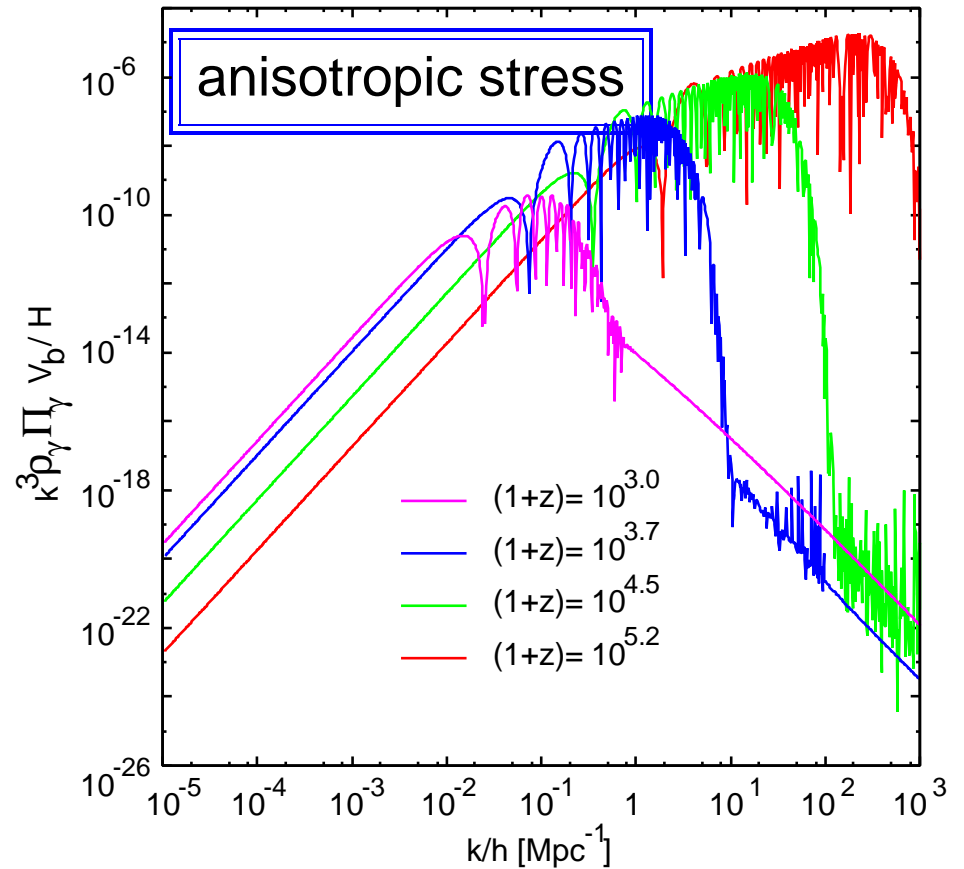
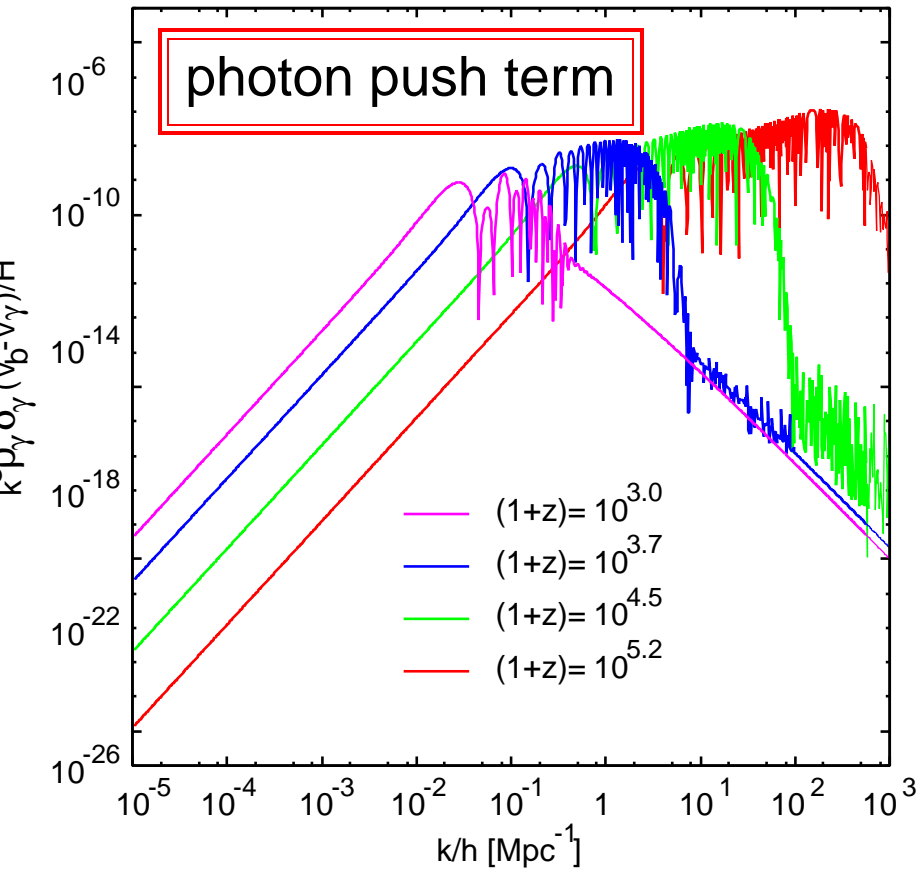
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source terms at various redshifts

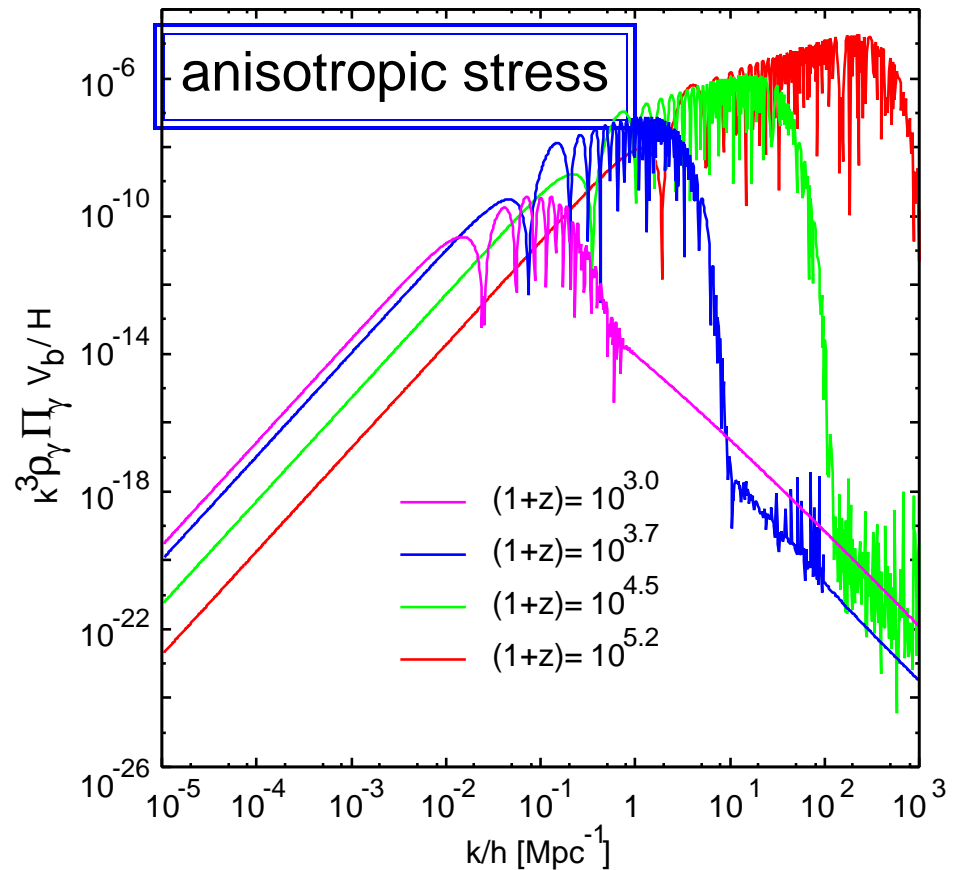
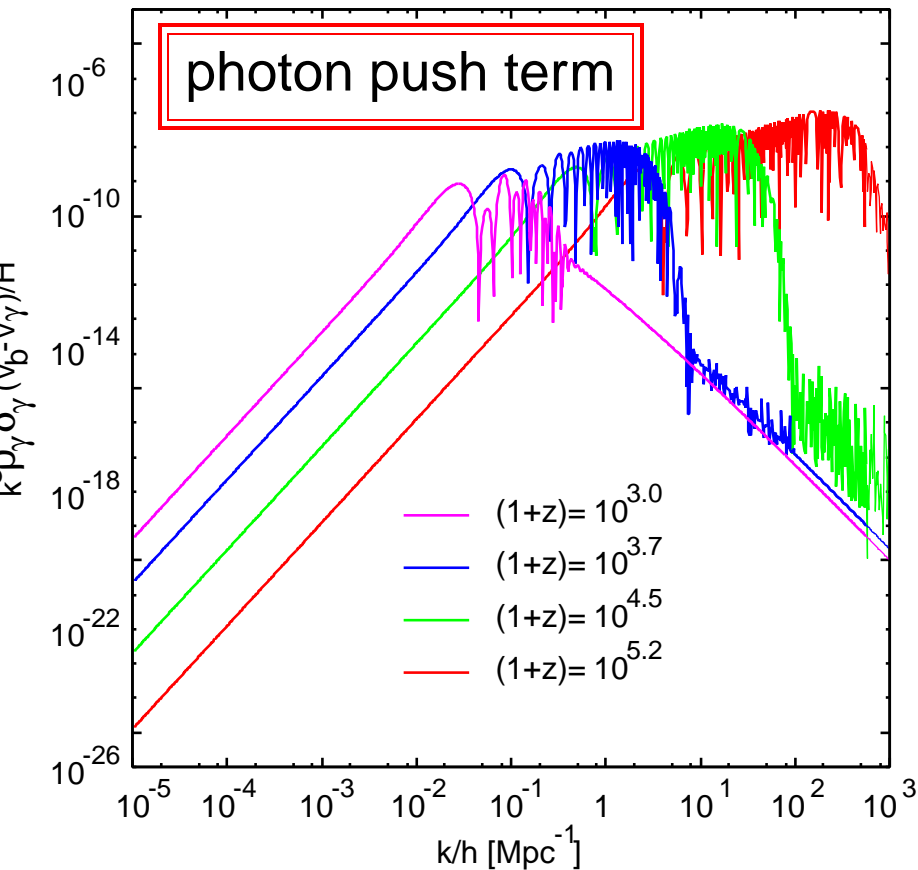
magnitudes of source terms per $\log z$. Magnetic fields can be obtained by time integration and non-linear convolution of these spectra.



$$\dot{B}^i = \frac{4\sigma_T^{(0)} \rho_\gamma}{3e} \epsilon^{ijk} \left[\delta_{\gamma,k}^{(1)} \left(u_{ej}^{(1)} - u_{\gamma j}^{(1)} \right) + \frac{1}{8} \left(u_{el,k}^{(1)} \Pi_{\gamma j}^{(1)} + u_{el}^{(1)} \Pi_{\gamma j,k}^{(1)} \right) \right],$$

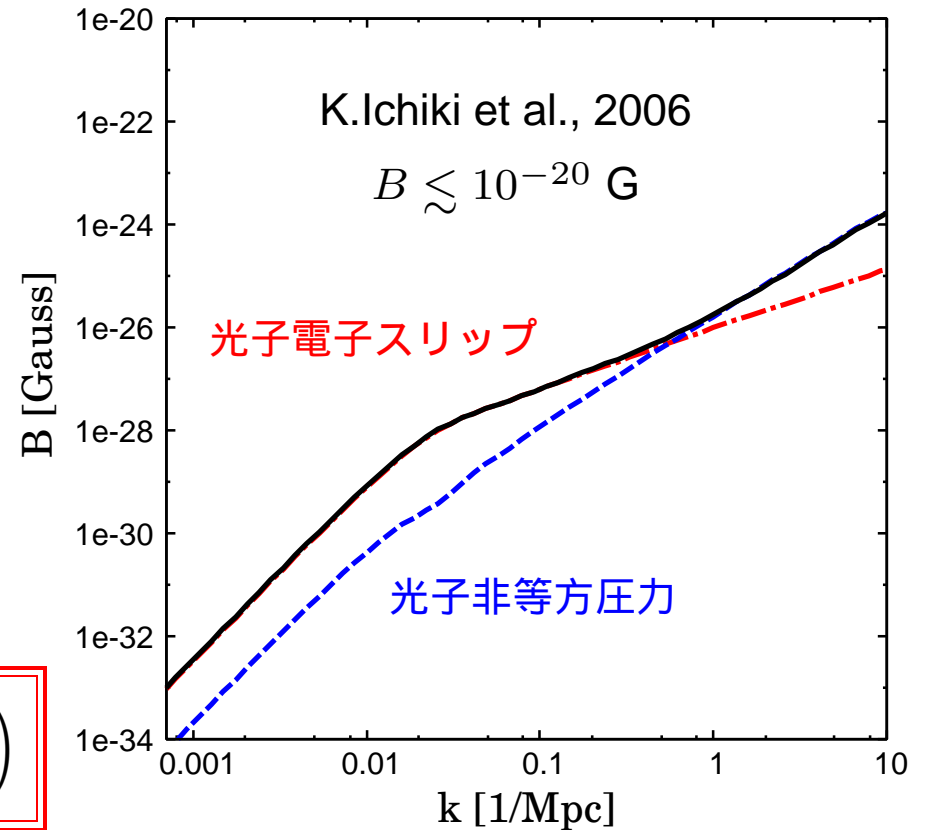
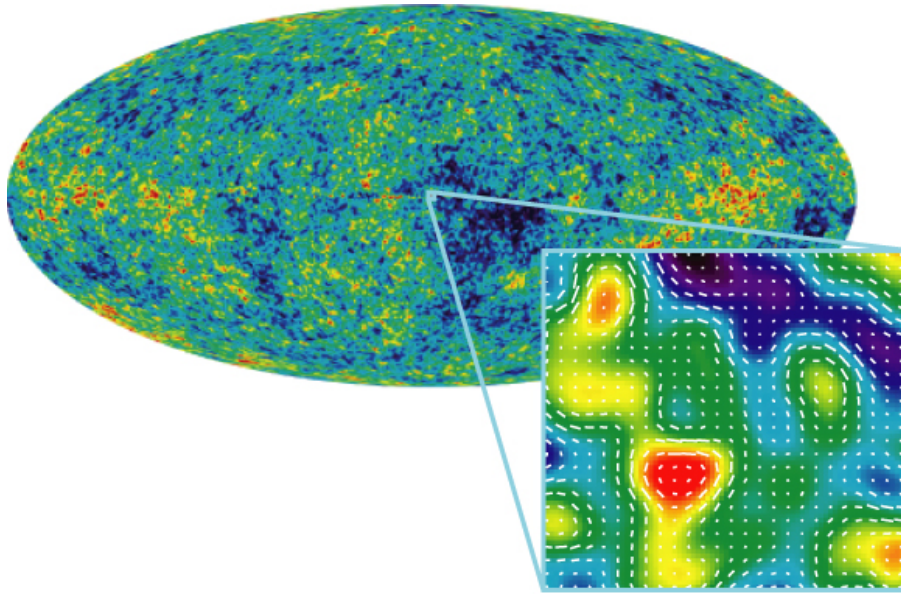
source terms at various redshifts

magnitudes of source terms per $\log z$. Magnetic fields can be obtained by time integration and non-linear convolution of these spectra.



B fields at the smaller scales are generated at the earlier epochs

Cosmological Magnetic Fields from Density Perturbations



$$\dot{B}^i = \frac{4\sigma_T}{3e} \rho_\gamma^{(0)} \epsilon^{ijk} \left[\delta_{\gamma,k}^{(1)} \left(u_{ej}^{(1)} - u_{\gamma j}^{(1)} \right) + \left(u_{ej,k}^{(2)} - u_{\gamma j,k}^{(2)} \right) + \frac{1}{8} \left(u_{el,k}^{(1)} \Pi_{\gamma j}^{(1)} + u_{el}^{(1)} \Pi_{\gamma j,k}^{(1)} \right) \right],$$

K. Takahashi et al., (2005)

Brief Summary

magnetic field generation from cosmological perturbations

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magnetic field generation from cosmological perturbations

- density fluctuation \rightarrow photon wind $\rightarrow \vec{j} \rightarrow \vec{B}$
- generated from second order effect, no assumptions (unavoidable generation)
- fields are correlated with CMB and Large scale structure

Brief Summary

magnetic field generation from cosmological perturbations

- density fluctuation \rightarrow photon wind $\rightarrow \vec{j} \rightarrow \vec{B}$
- generated from second order effect, no assumptions (unavoidable generation)
- fields are correlated with CMB and Large scale structure
- characteristic values
 - 10^{-20} G @ 10 kpc
 \rightarrow sufficient for galactic dynamo
 - 10^{-24} G @ 1 Mpc
 \rightarrow strict lower bound at intergalactic fields

An Interesting Speculation

Inflation
generation of density perturbations

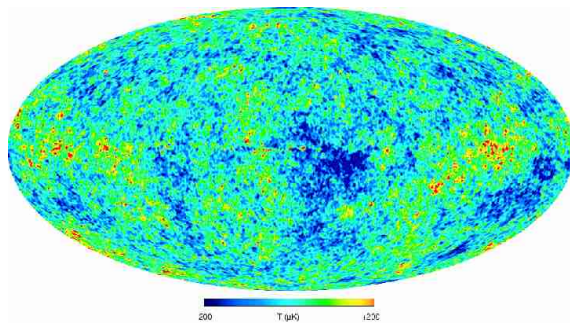


$$ds^2 = (g_{\mu\nu} + h_{\mu\nu})dx^\mu dx^\nu$$
$$G_{\mu\nu} = 8\pi GT_{\mu\nu}$$

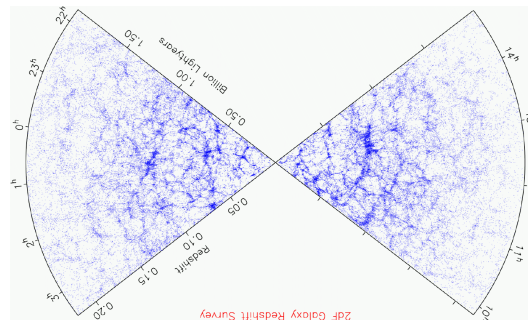
cosmological density perturbation theory



CMB anisotropies



Large Scale Structure



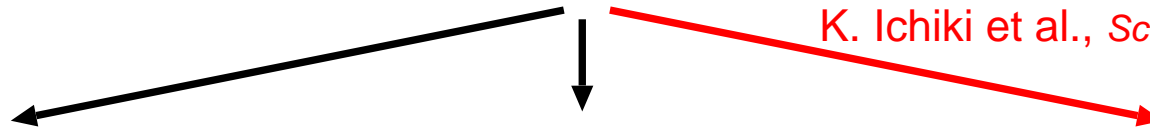
An Interesting Speculation

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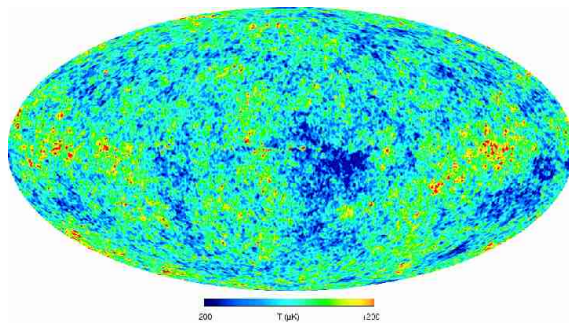
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cosmological density perturbation theory

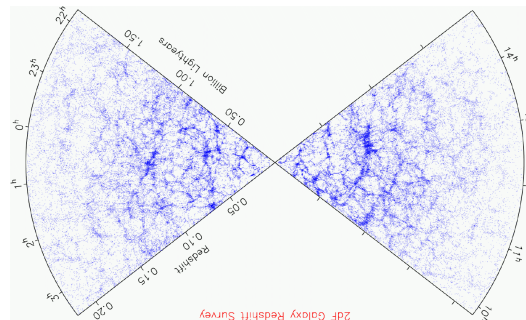


K. Ichiki et al., *Science*, (2006)

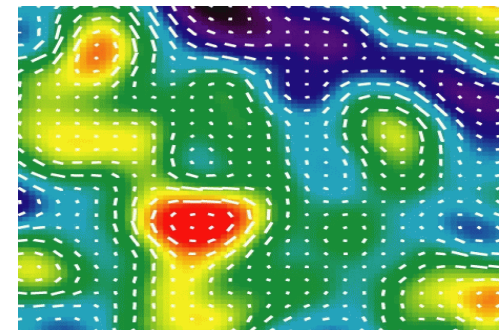
CMB anisotropies



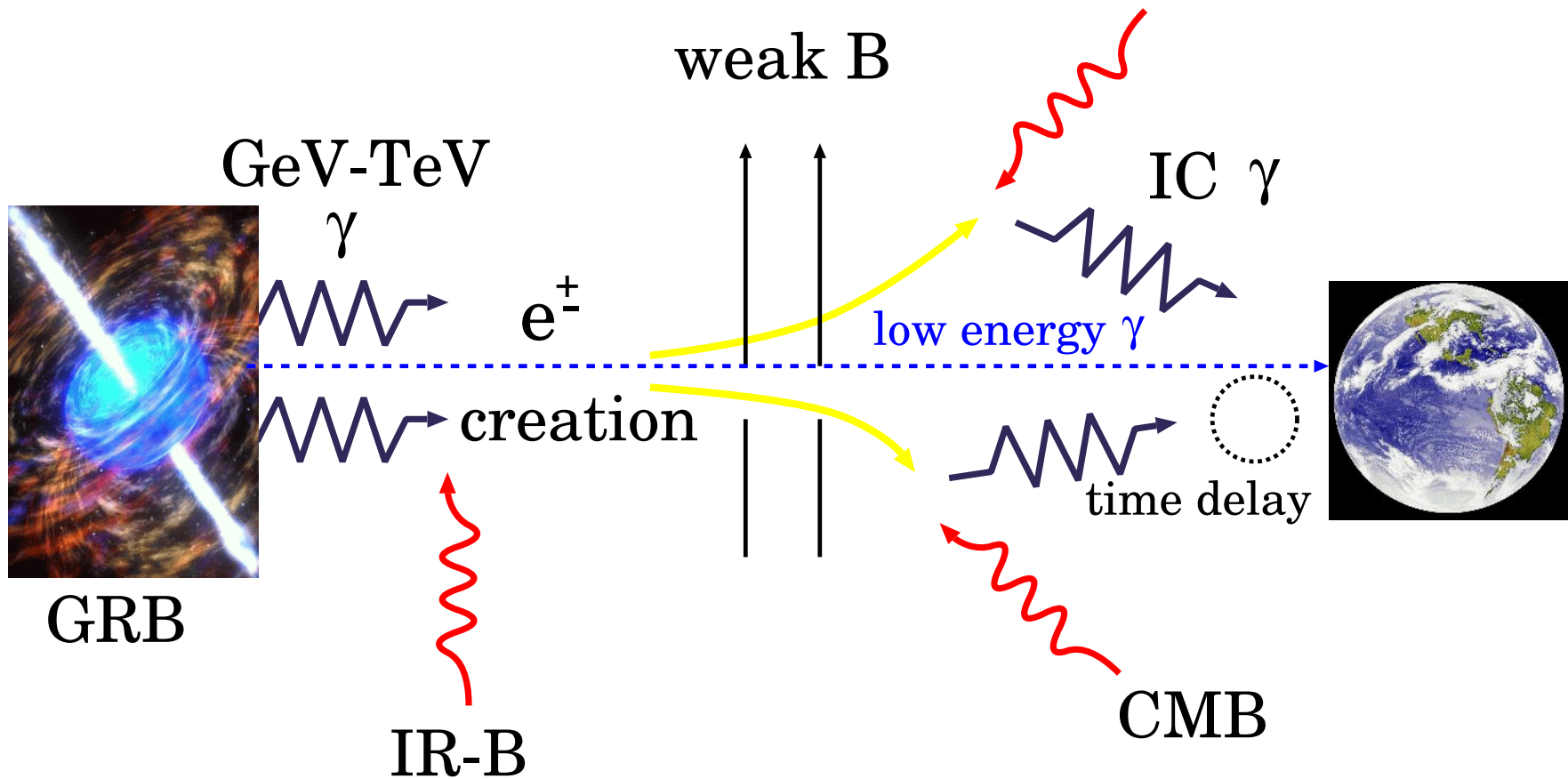
Large Scale Structure



Cosmic B fields



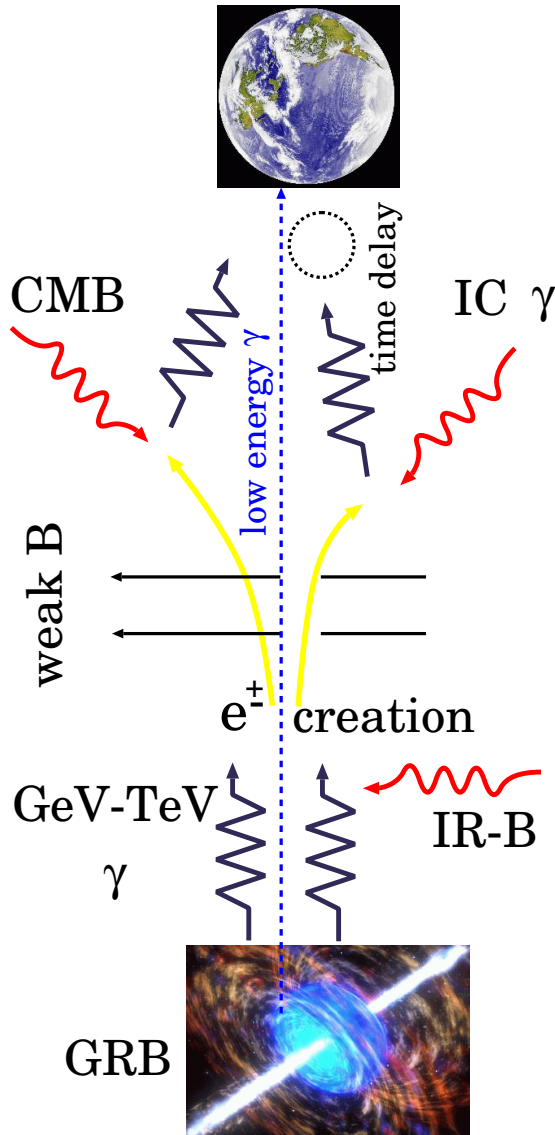
Detectability (R. Plaga, *Nature* 1995)



- arrival time should be delayed in the presence of intergalactic magnetic fields.

Our model to predict delayed emission

improved version of Razzaque et al., ApJ (2004)



$$\frac{d^2 N_{\text{delay}}^{\text{IC}}}{dt dE_\gamma} = \int d\gamma_e \frac{dN_e}{d\gamma_e dt} t_{\text{GRB}} \frac{d^2 N_\gamma^{\text{IC}}}{dt dE_\gamma} \left(\frac{t_{\text{IC}}}{\Delta t} \right) e^{-\tau_{\gamma\gamma}(E_\gamma)}$$

↑
delayed Δt + re-emitting γ -rays by IC

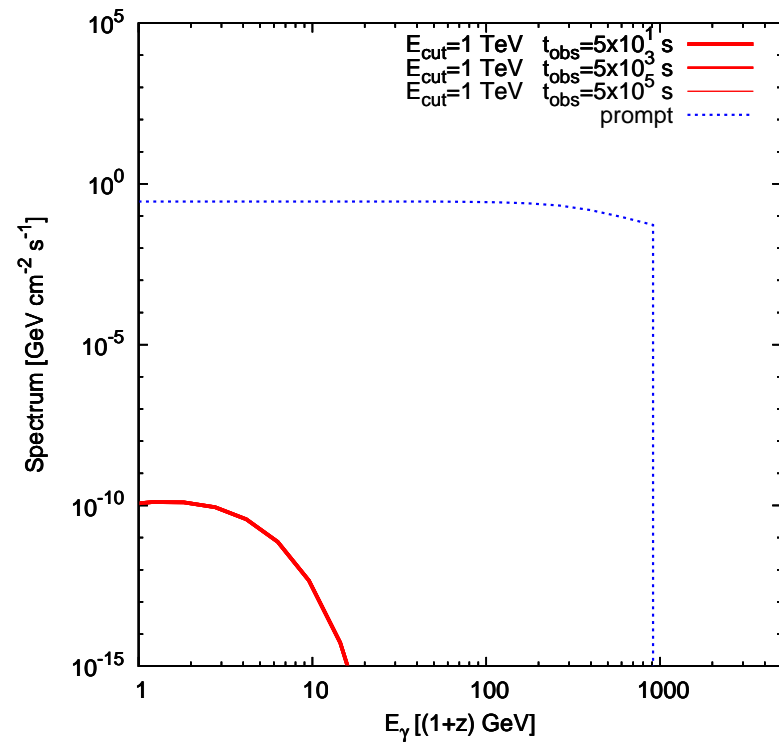
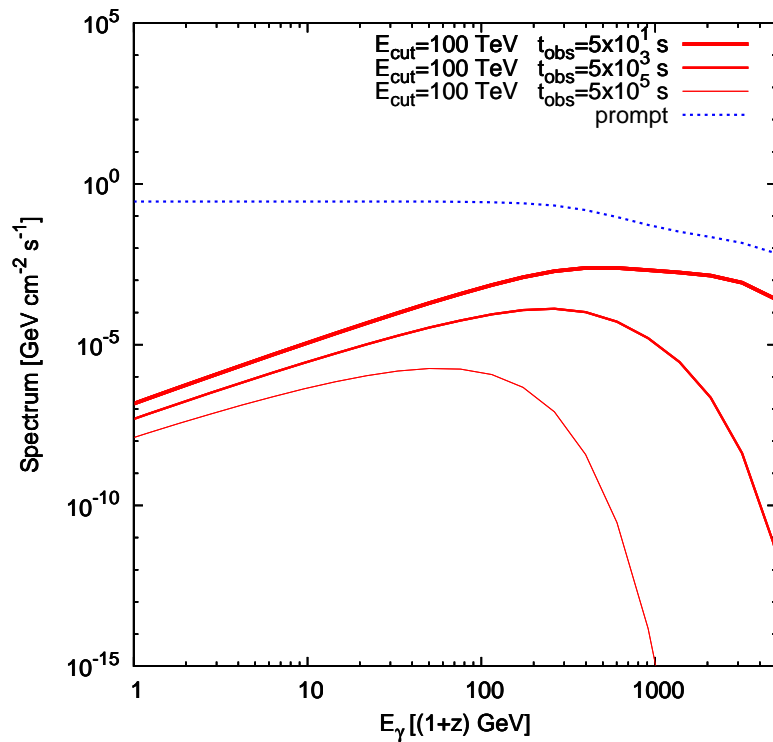
↑
 $\epsilon_\gamma \rightarrow 2m_e \gamma_e$ (e^\pm 生成), $\tau_{\gamma\gamma}$: optical depth of IRB

$$\frac{d^2 N_e}{2m_e d\gamma_e dt} = \frac{L_{\gamma, \text{iso}}}{4\pi D_L^2} \frac{\alpha-1}{\epsilon_{\gamma, \text{pk}}^2} \left(\frac{2m_e \gamma_e}{\epsilon_{\gamma, \text{pk}}} \right)^{-\alpha} (1 - e^{-\tau_{\gamma\gamma}(2m_e \gamma_e)})$$

↑
simplest power law (synchrotron) spectrum

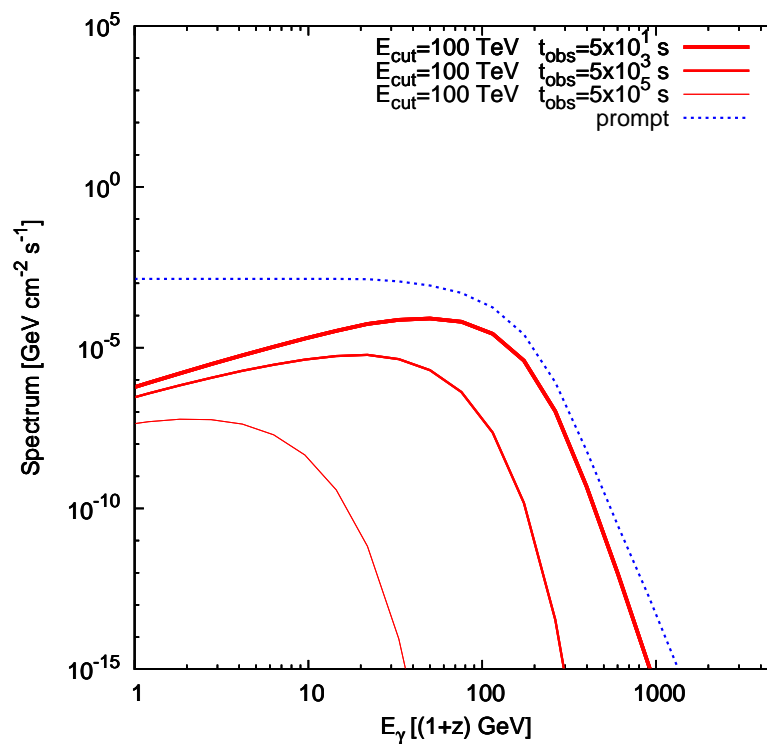
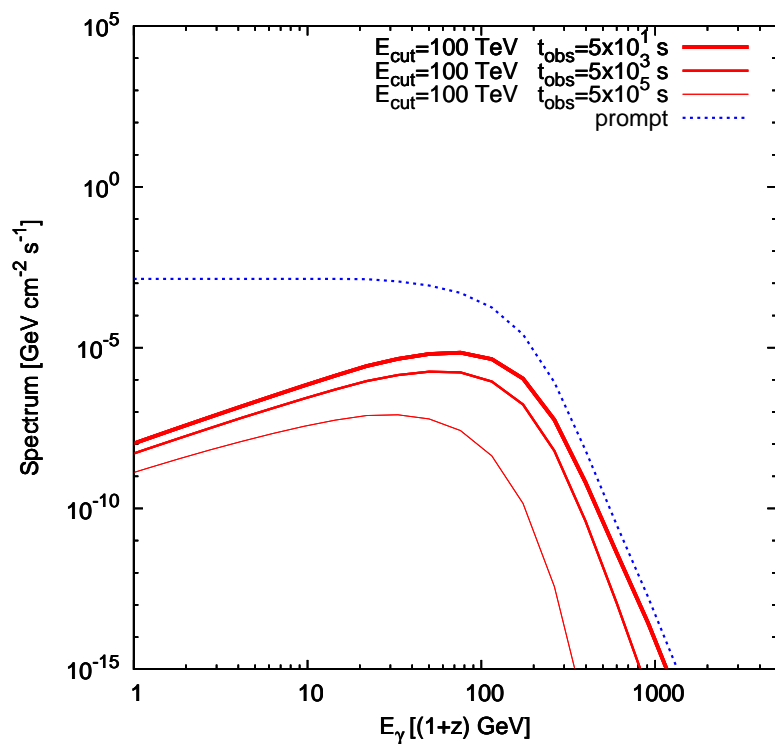
$$\frac{d^2 N_\gamma}{d\epsilon_\gamma dt} = \frac{L_{\gamma, \text{iso}}}{4\pi D_L^2} \frac{\alpha-1}{\epsilon_{\gamma, \text{pk}}^2} \left(\frac{\epsilon_\gamma}{\epsilon_{\gamma, \text{pk}}} \right)^{-\alpha} (\epsilon_\gamma > \epsilon_{\gamma, \text{pk}})$$

Delayed spectrum ($B=10^{-18}$ G; $z=0.1$)



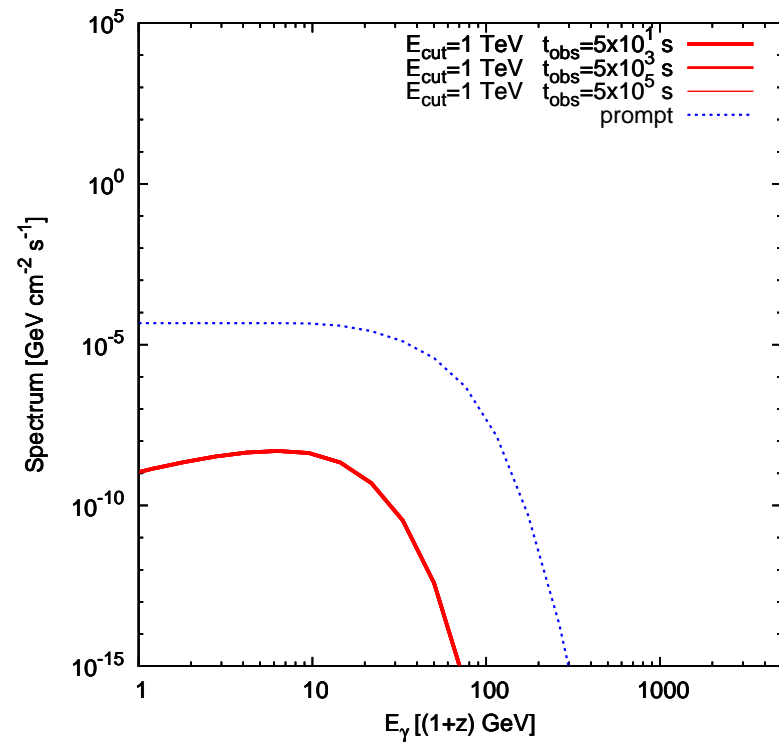
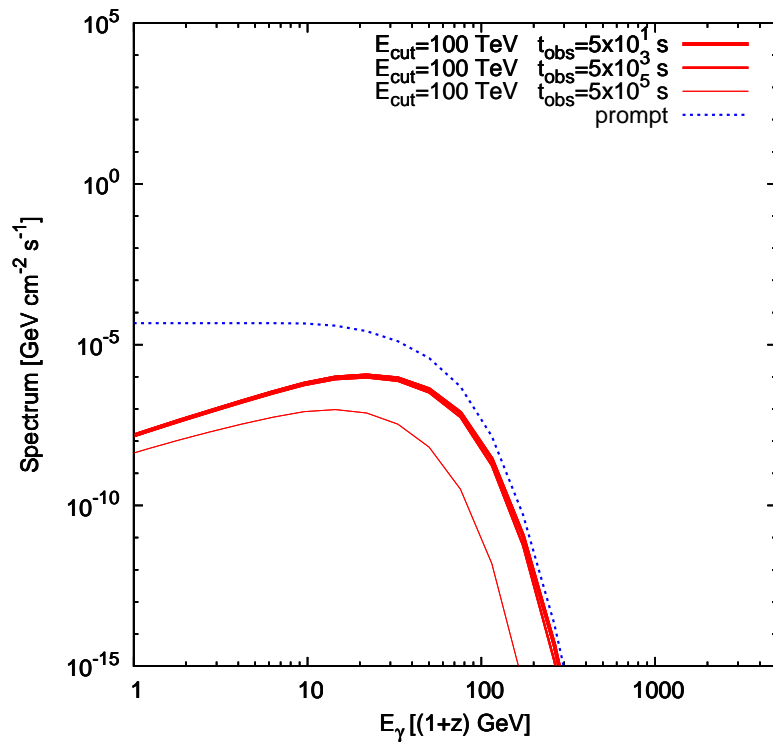
- IRB absorption will be important $E_{\gamma} > 100$ GeV, but not significant for γ -ray bust at $z = 0.1$.

Delayed spectrum ($B=10^{-18}$ G; $z=1.0$)



- Time dependence of the delayed spectra depends on the magnetic field configuration.
→ information on magnetic field spectrum

Delayed spectrum ($B=10^{-18}$ G; $z=4.0$)

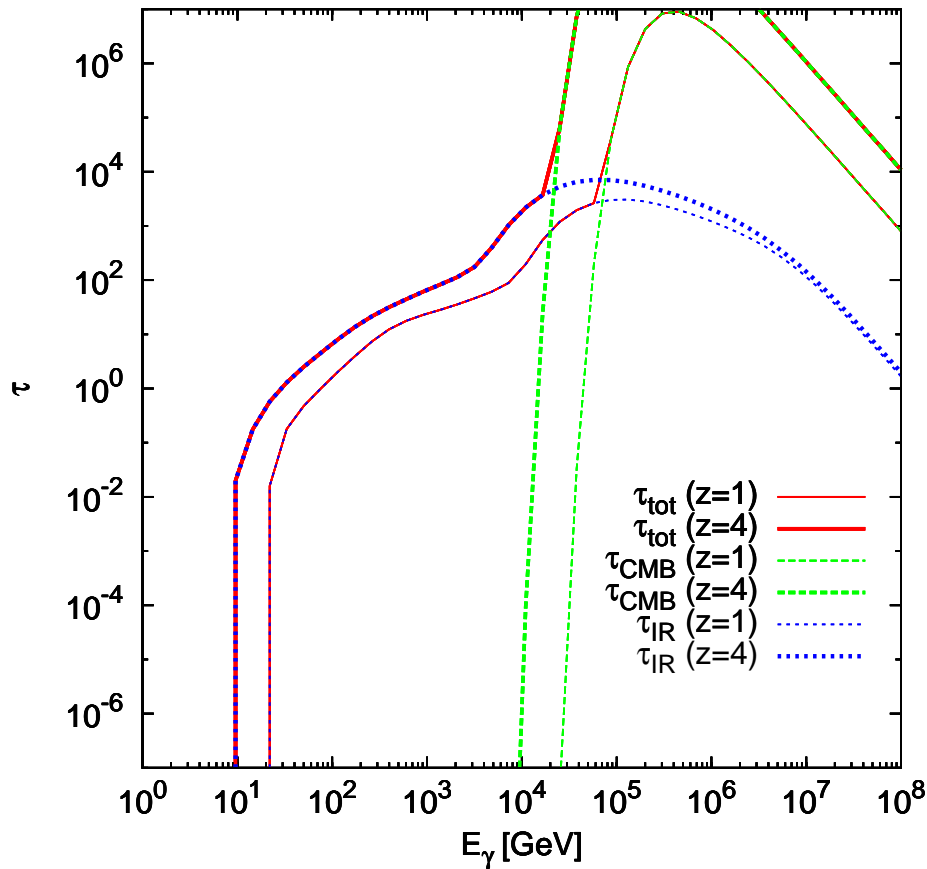


- Delayed IC-photons are absorbed once again for γ -ray bursts at $z = 4$.

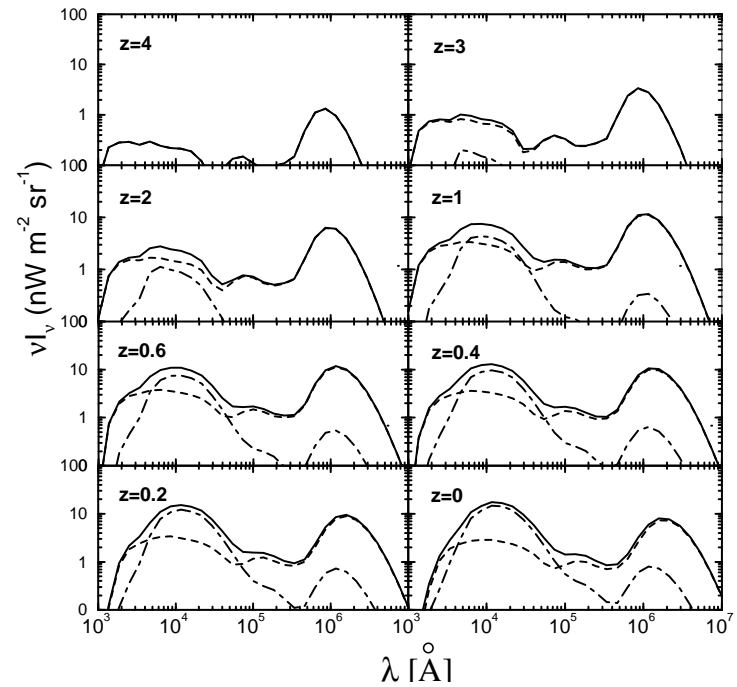
Summary

- The origin of large scale magnetic fields is still unknown
 - Cosmological ? (predictable (but model dependent))
 - Astrophysical ? (well known physics, but highly uncertain)
 - We showed that cosmological density fluctuations naturally generate magnetic fields (although very weak).
 - characteristic values
 - 10^{-20} G @ 10 kpc
→ sufficient for galactic dynamo
 - 10^{-24} G @ 1 Mpc
→ strict lower bound at intergalactic fields
- Magnetic fields in void regions can be probed by GeV-TeV delayed photons from γ -ray bursters
 - Fields as weak as $B \sim 10^{-18}$ G may be detectable.
 - Amplification mechanisms are needed for cosmologically generated magnetic fields to be detectable and/or dynamically important.

optical depth for $\gamma + \gamma \rightarrow e^\pm$



- Model employed for IRB
Kneiske et al., A&A, 386, 1, (2002)



important scales involved

Can we probe magnetic fields in void regions?

- energy of electrons produced by γ -rays with $E_\gamma = 20\text{TeV}$:

- $E_e = 10.0\text{TeV} \left(\frac{E_\gamma}{20\text{TeV}} \right)$, $\gamma_e = 2 \times 10^7 \left(\frac{E_\gamma}{20\text{TeV}} \right)$

- $\lambda_{\gamma\gamma} \sim \frac{1}{\sigma_T n_{IR}} = 6\text{Mpc} \left(\frac{n_{IR}}{1\text{cm}^{-3}} \right)$ (\leftarrow creation time)

- e^\pm creation in void

- How long distance can this e^\pm propagate?

- $\lambda_{IC} = \frac{3m_e}{4\gamma_e\sigma_T u_{\text{cmb}}}$
 $= 100\text{kpc} \left(\frac{E_\gamma}{20\text{TeV}} \right)^{-1}$

- cool down in void

