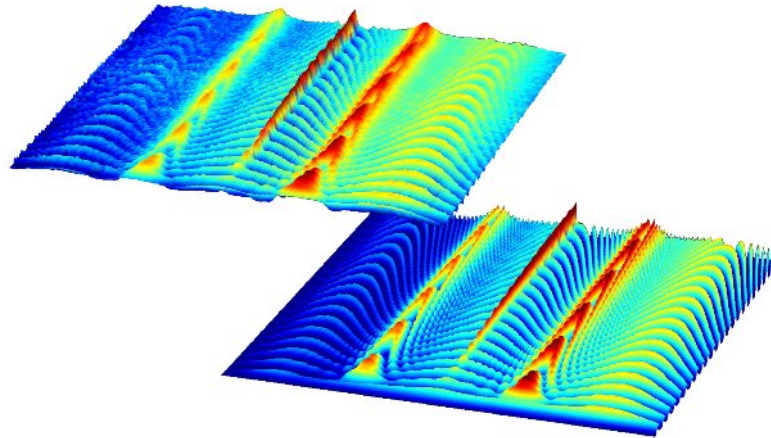


Controlled Multiple Quantum Coherences of Nuclear Spins in a Nanoscale Device



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Y. Hirayama (NTT BRL, SORST-JST)

Quantum Computers using Coherence of Nuclear Spins

Quantum computation based on NMR

Experimental realization of D-J algorithm

I. L. Chuang *et al.*, Nature (1998).

Experimental realization of Shor's factoring algorithm

L. M. Vandersypen *et al.*, Nature (2001).

PROBLEM

Liquid-state NMR is not suitable
for a scalable device

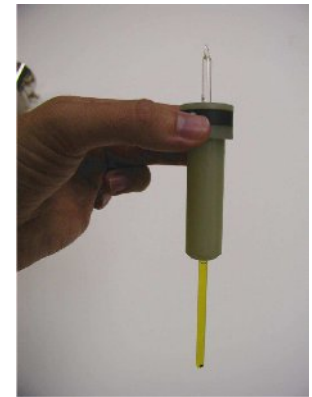


Quantum computers based on
all-electrical solid-state NMR

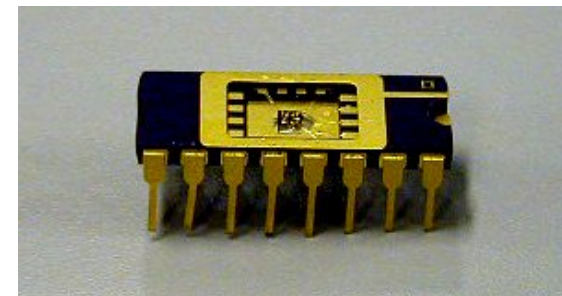
Electronics + NMR

A typical NMR sample

L. M. K. Vandersypen
quant-ph/0205193v1

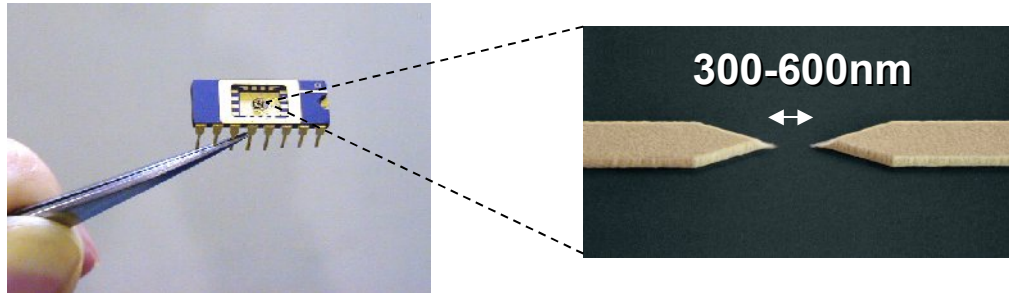


Our NMR chip



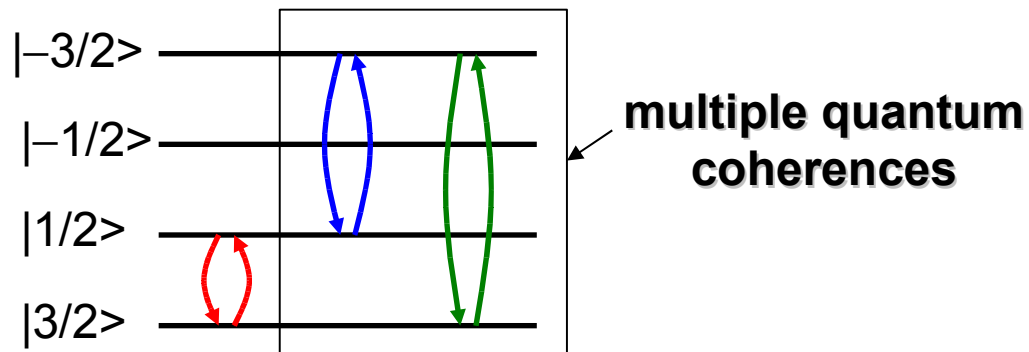
Outline

A self-contained semiconductor chip that can access nuclear spins in a nanoscale region



New NMR: Direct detection of multiple quantum coherences, which are invisible by conventional NMR

Nuclei often possess total spin angular momentum greater than a half and multiple spin levels are formed.

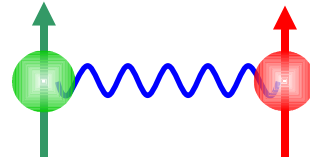


Control and Detection of Coherence of Nuclear Spins

Contact Hyperfine Interaction

Nuclear Spins

Long coherence time

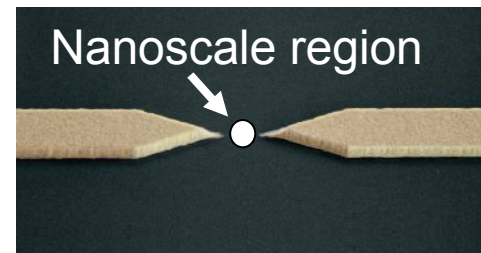


Electron Spins

Electrical accessibility

① Initialization (Polarization)

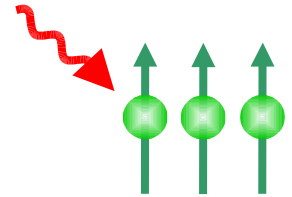
- Electron-nuclear spin coupling (flip-flop process).
->Polarization only in the **nanoscale region**.



② Creation of quantum mechanical superpositions of states

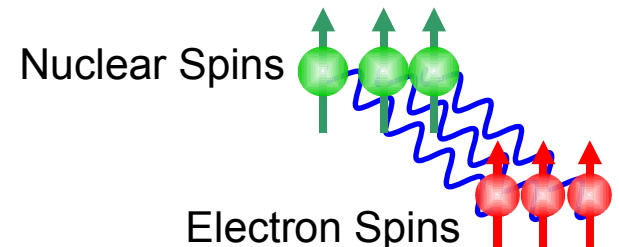
- **Nuclear magnetic resonance (NMR)** on a self-contained **semiconductor chip**

r.f. field



③ Readout

- Electron-nuclear spin coupling.
- Measurement of polarization by resistance



Polarization of Nuclear Spins

Electron-nuclear spin coupling in 2D system
in the **fractional quantum Hall regime** ($\nu=2/3$)

- **Long-time-scale resistance enhancement (~ 10 min)**

- ➡ Due to long longitudinal (spin-lattice) relaxation time T_1

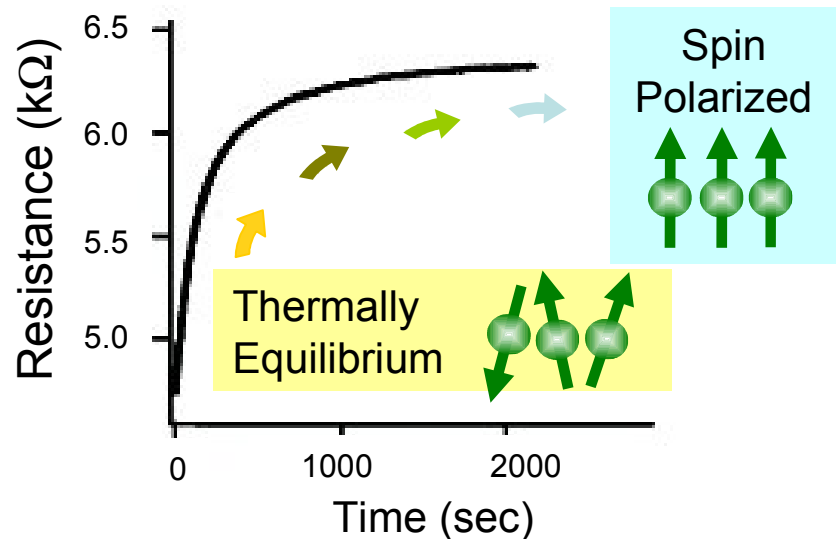
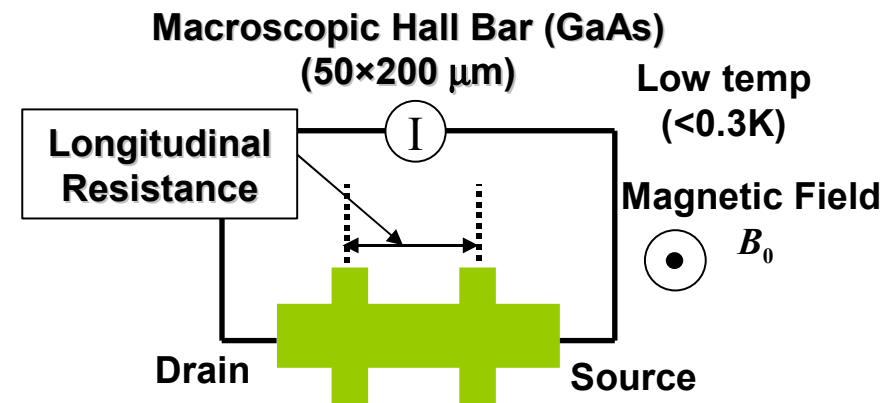
- **Current dependence**

- ➡ Large current is required to polarize nuclear spins

- **Resistively detected CW-NMR**

- Kronmuller *et al.*, PRL **81**, 2526 (1998).
- Smet *et al.*, PRL **86**, 2412 (2001).
- Hashimoto *et al.*, PRL **88**, 176601 (2002).
- Smet *et al.*, Nature **415**, 281 (2002).
- Kumada *et al.*, PRL **89**, 116802 (2002).
- Kraus *et al.*, PRL **89**, 266801 (2002).

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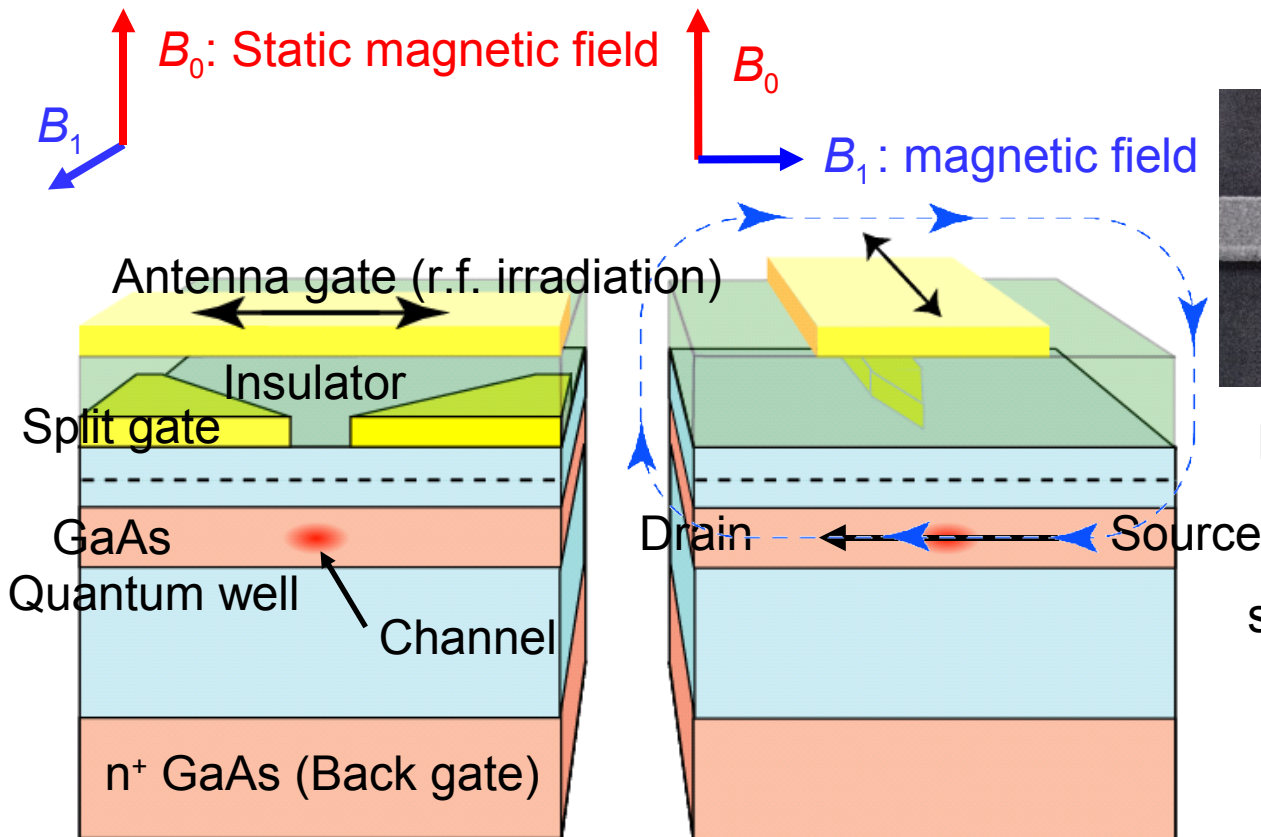


Hashimoto *et al.*, PRL (2002).

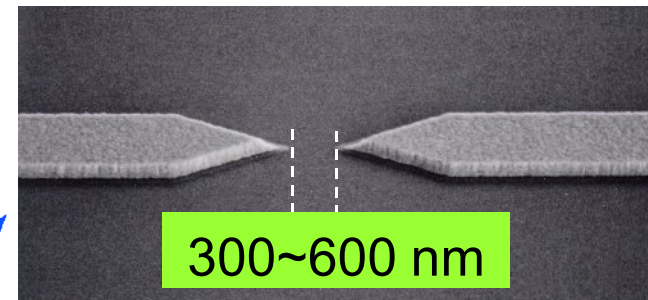
Self-contained NMR Device

Three gate electrodes

1. **Back gate**: to control electron density.
2. **Split gate**: to **define nanoscale region**
3. **Antenna gate**: for **r.f. irradiation**



SEM Image

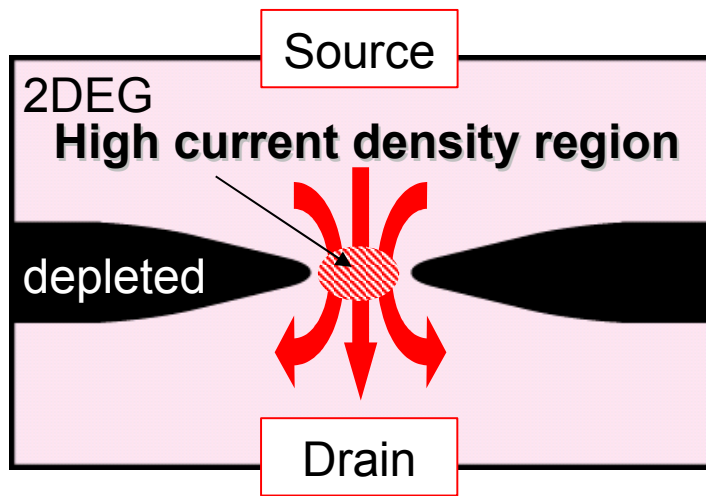


Number of Nuclei $< \sim 10^8$
This is much smaller than
the **detection limit** of
standard NMR ($\sim 10^{11}-10^{13}$).

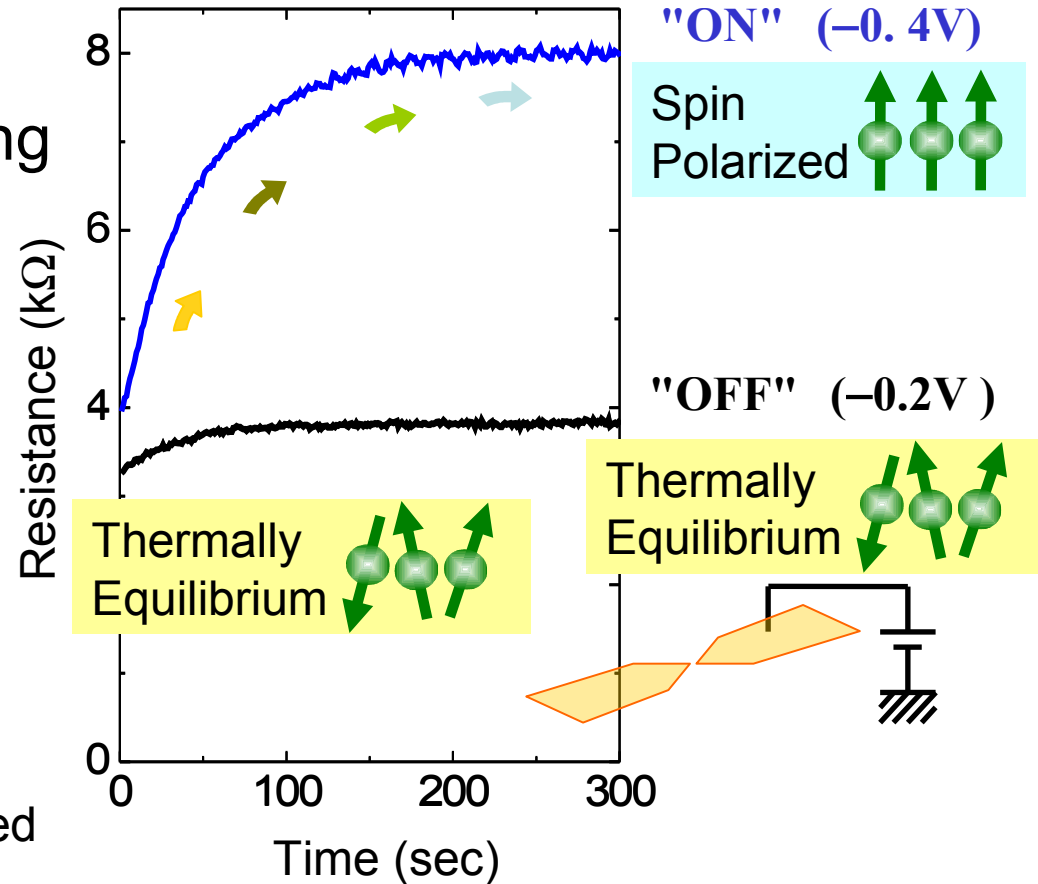
Polarization of Nuclear spins in a Nanoscale Region

Temperature ~ 0.1 K

Current dependence of electron-nuclear spin coupling

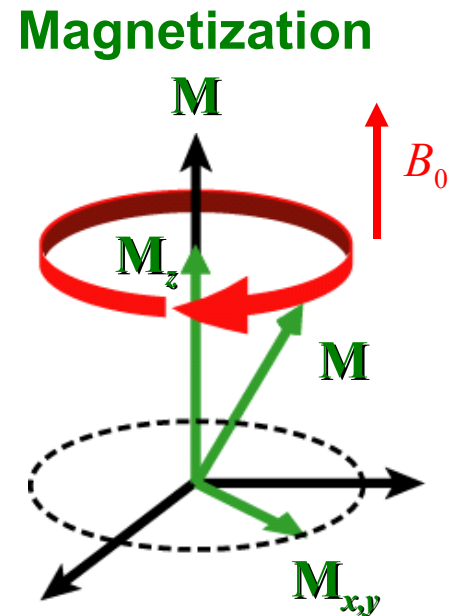
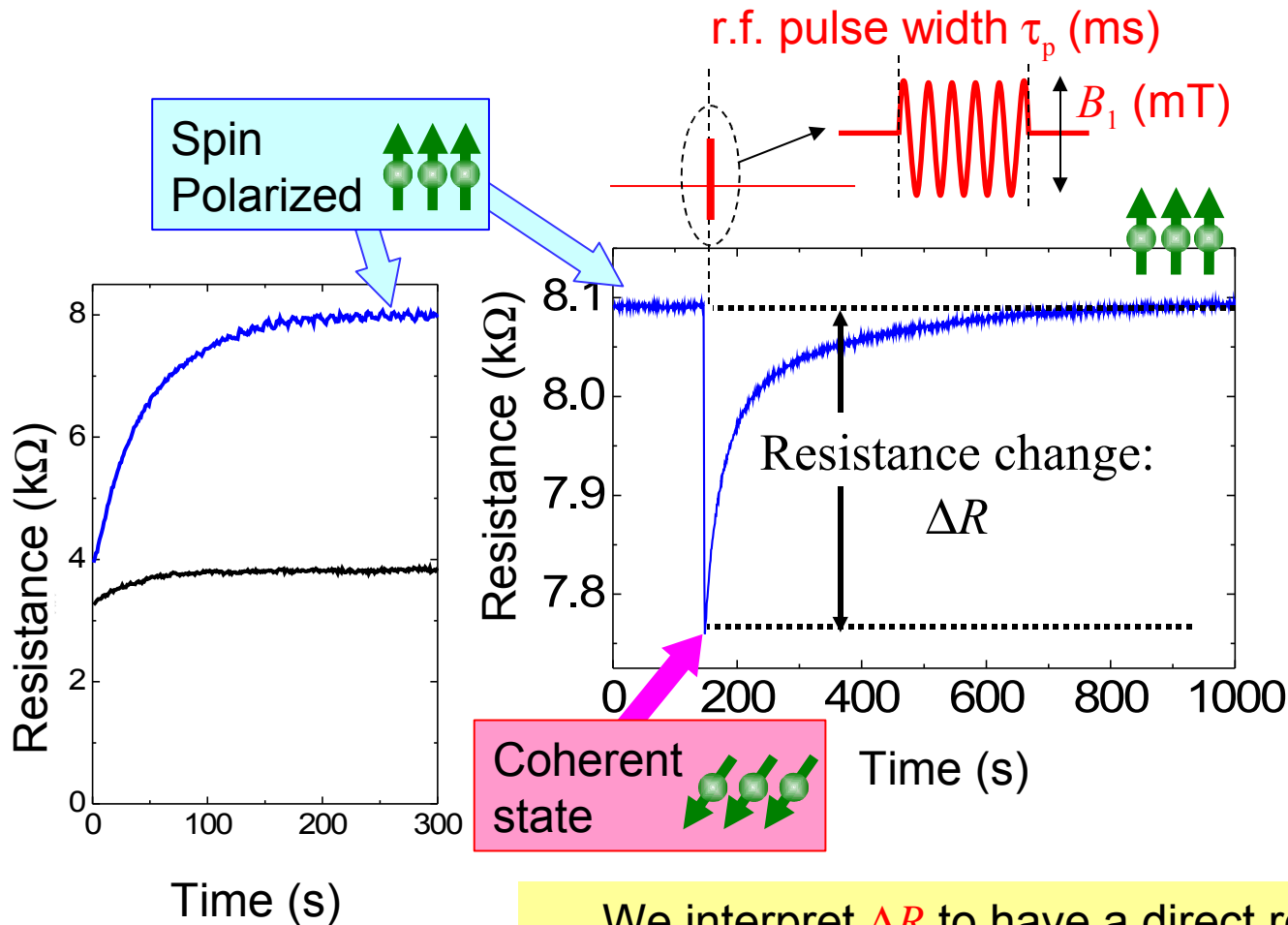


By producing high current density region, nuclear spins can be polarized only in the nanoscale region.



r.f.-Pulse Irradiation

to create a quantum mechanical superposition of states

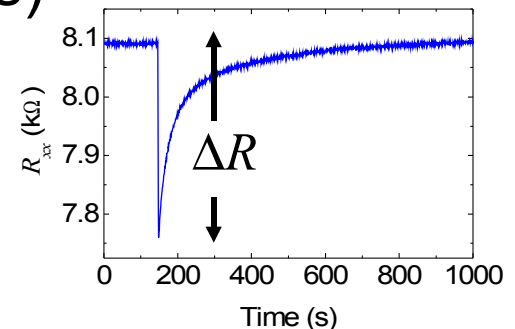
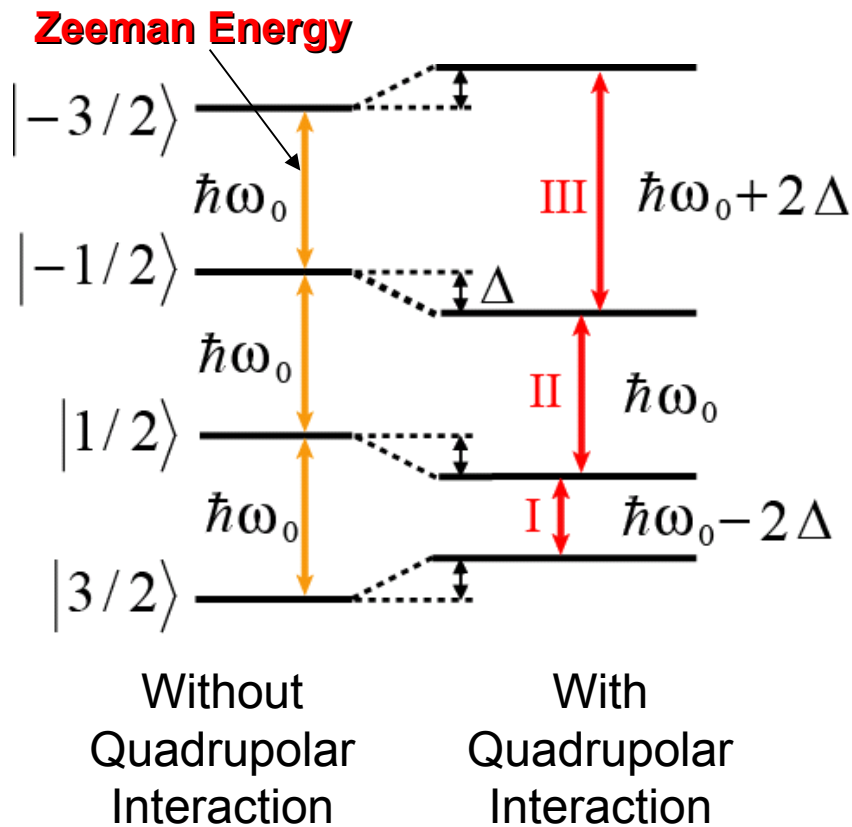


We interpret ΔR to have a direct relationship with the change in the **z component of the magnetization, M_z** induced by the altered population by the r.f. pulse.

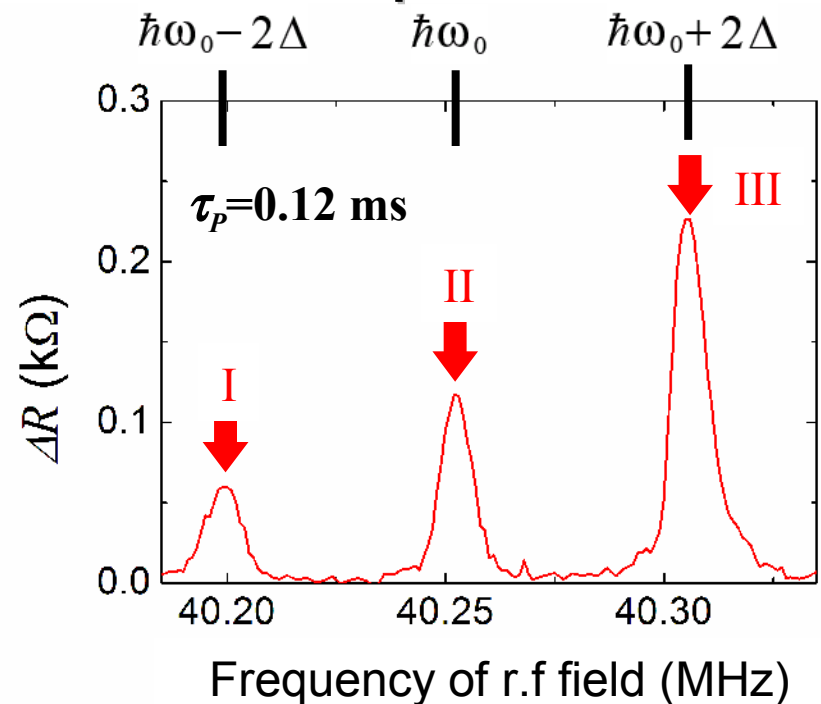
Quadrupolar Interaction

Nuclear spin $I=3/2$ (^{69}Ga , ^{71}Ga , ^{75}As)

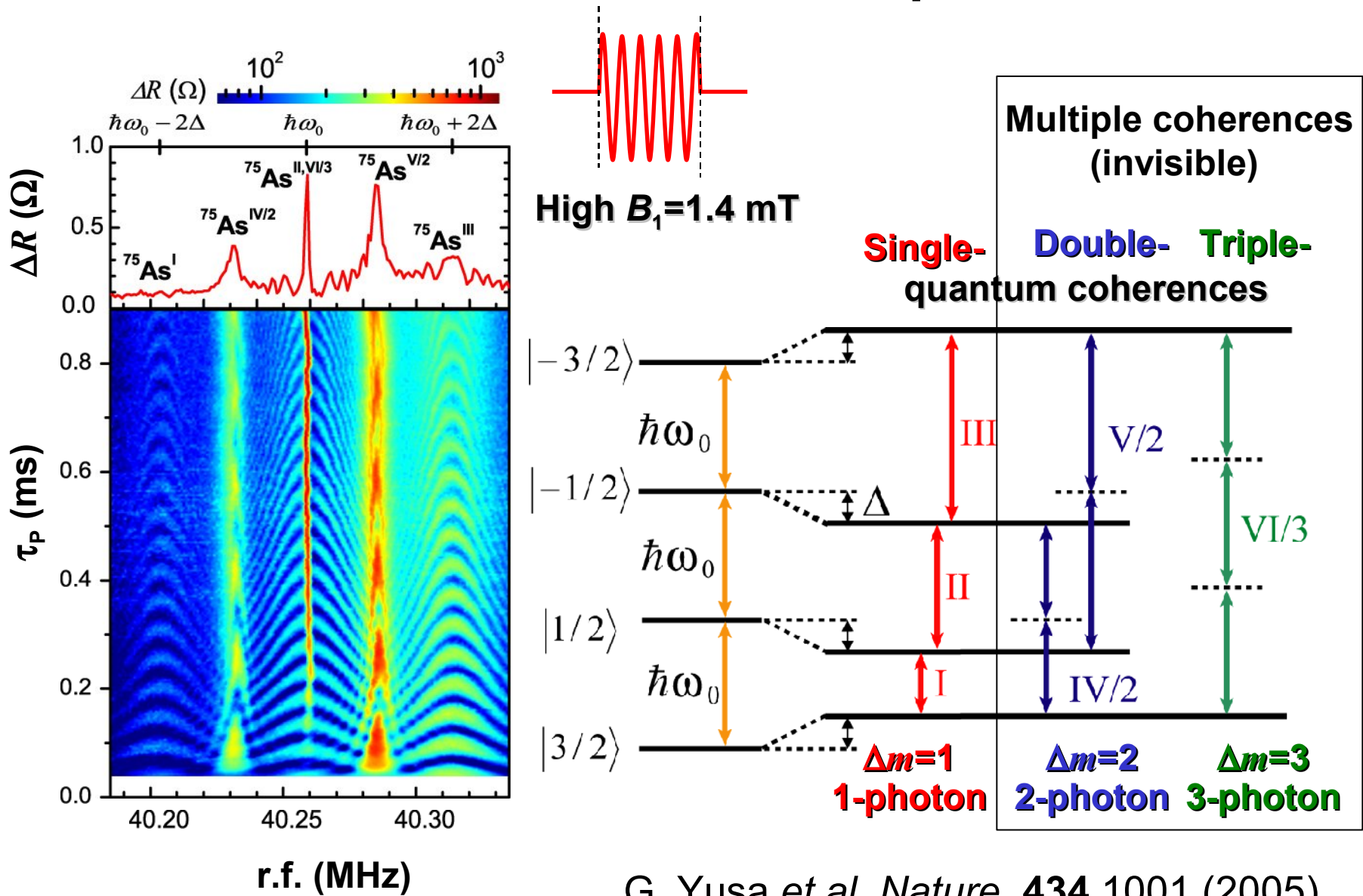
Energy Diagram of $I=3/2$ system



NMR Spectrum of ^{75}As



Time Evolutional Spectra



Quantitative Discussion

ASSUMPTION:

The change in the resistance (ΔR) is **proportional** to the change in the magnetization M_z induced by altered population: $\Delta R \propto \Delta M_z$

Simulation of M_z using **rotating-frame approximation**

$$i\hbar \frac{d}{dt} |\Psi(t)\rangle = H |\Psi(t)\rangle$$

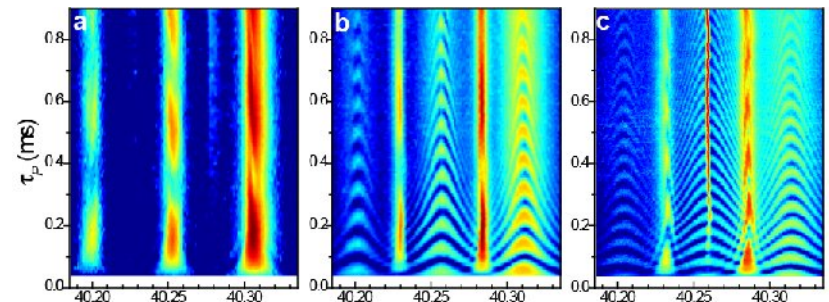
$$H = \underbrace{-\gamma\hbar B_0 I_z}_{\text{Zeeman}} + \underbrace{\Delta / 3 [3I_z^2 - I(I+1)]}_{\text{Quadrupolar}} + \underbrace{\gamma\hbar B_1 \cos(\omega t) I_x}_{\text{r.f. field}}$$

γ : gyromagnetic ratio
 Δ : quadrupolar constant

The initial population is the same for all three panels.

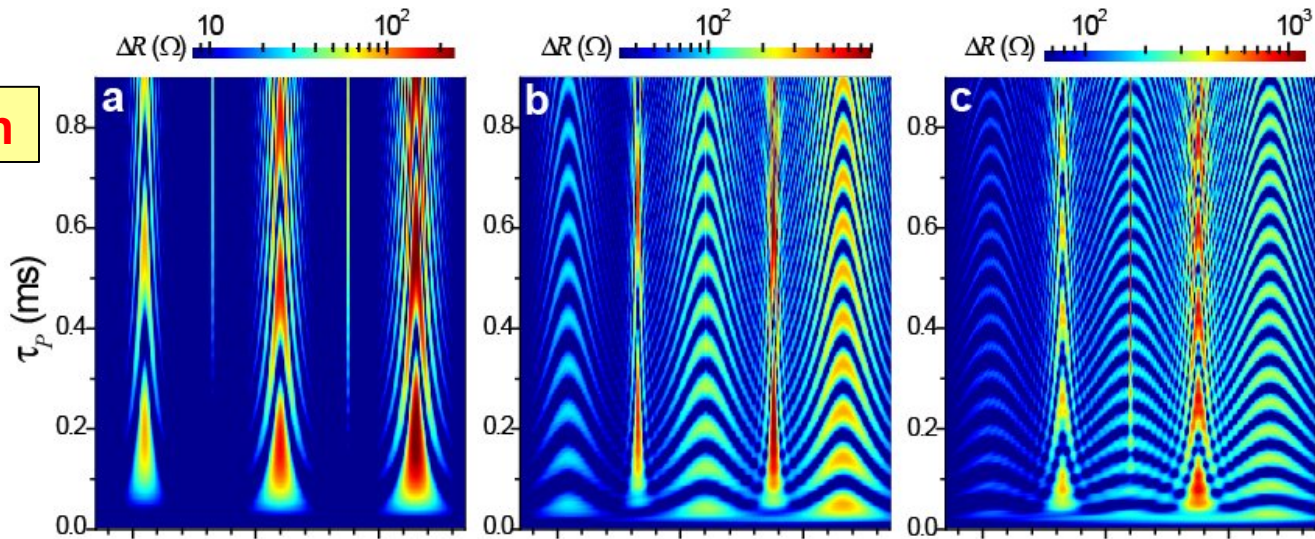
B_1 is the only **input parameter** for each panel.

Experiment

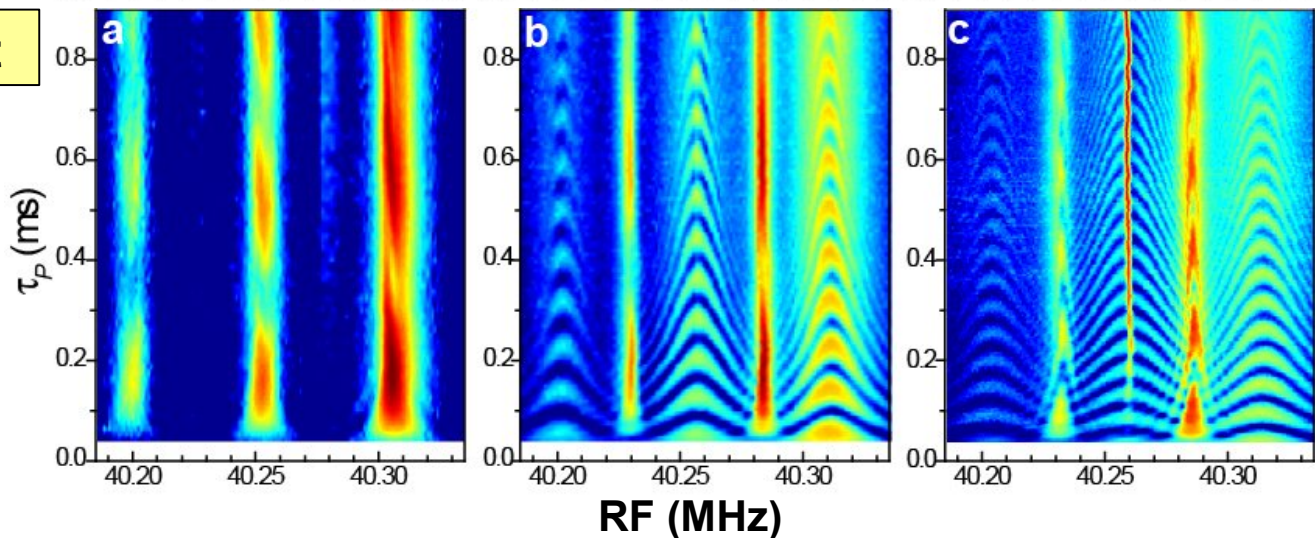


Numerical Calculations

Calculation



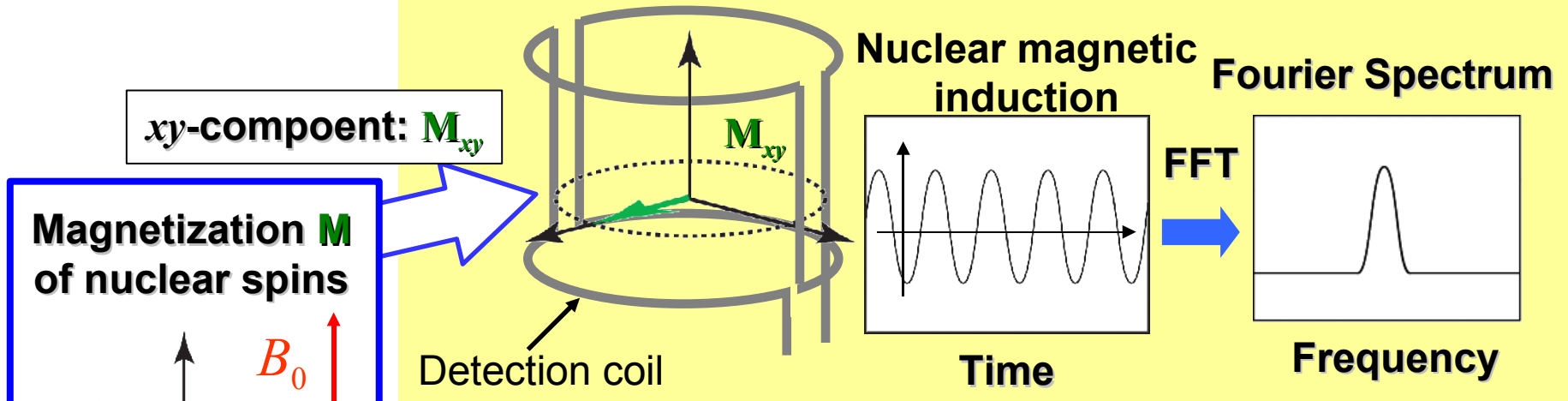
Experiment



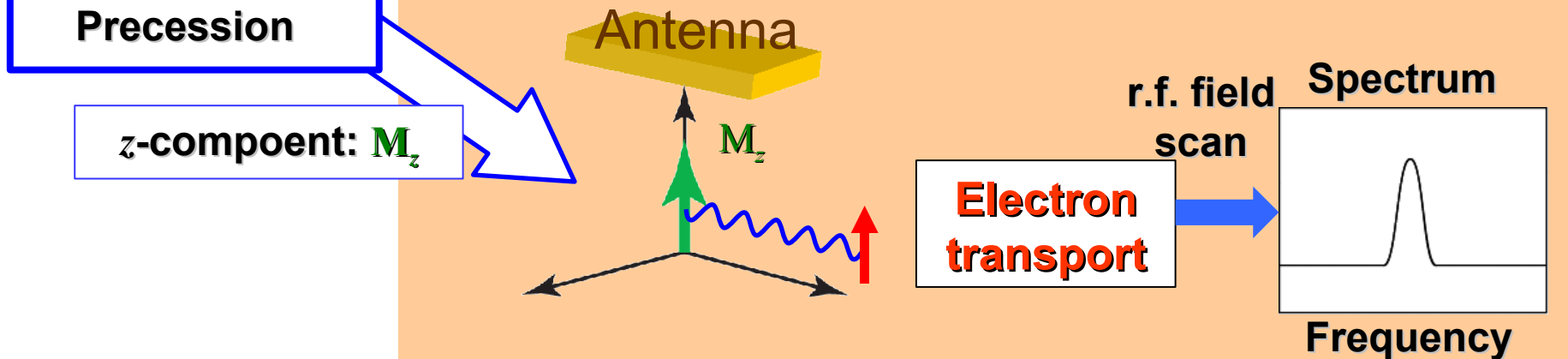
We measure the change in M_z by the resistance

Standard vs. Our NMR

Standard NMR (induction detected)



Our NMR

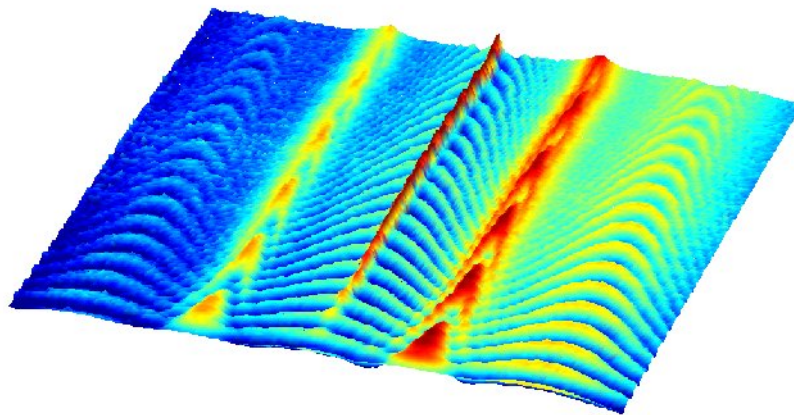


Our NMR (M_z detection)

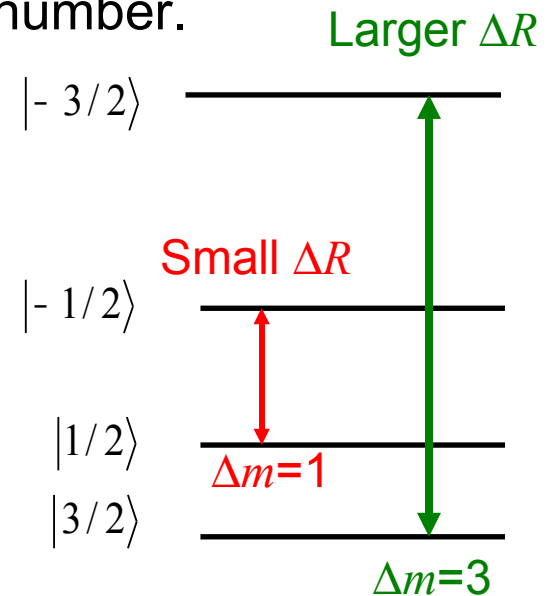
- **Direct detection** of multiple quantum coherence

- The oscillation amplitudes for **higher order coherences** are **larger** than those for single, reflecting **greater change in spin**

- Width is scaled down by the photon number.



on a logarithmic scale

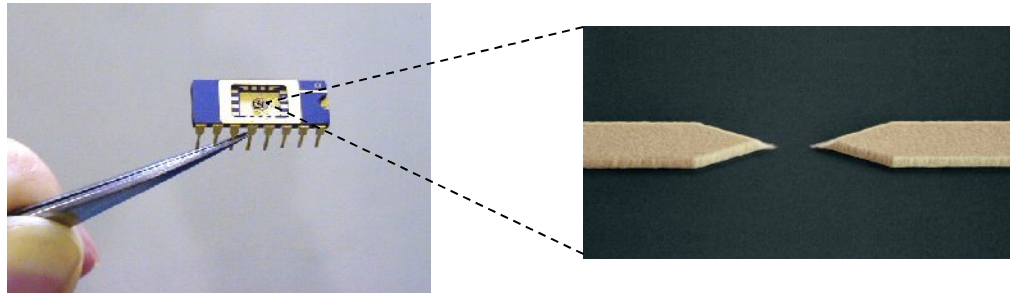


- **High sensitivity** ($<10^8$)

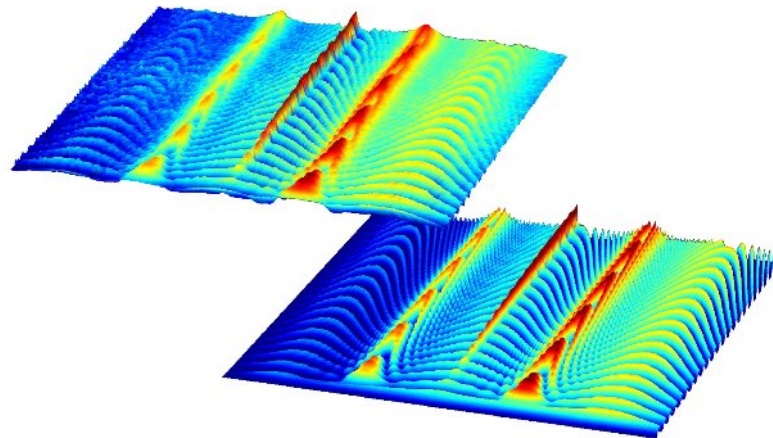
- Much less than the detection limit of standard NMR (10^{11} - 10^{13})

Conclusion

A self-contained semiconductor device with a **novel paradigm of NMR**, which accesses nuclear spins in a nanoscale region



Direct detection of **multiple quantum coherences**, which is **invisible** by conventional NMR



Three single, **two double** and **one triple** quantum coherences for one nuclide, 18 in total for three nuclides (^{69}Ga , ^{71}Ga , ^{75}As), are completely controlled by all-electrical means.