



# Designing Molecular Swimmers

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The 21<sup>st</sup> Century Center-of-Excellence Program  
Exploring New Science by Bridging Particle-Matter Hierarchy  
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# Swimming and Length Scale

- Simplified Dynamics

$$m \frac{d^2 x}{dt^2} + \mathbf{V} \frac{dx}{dt} = F(t)$$

$F(t)$  is Periodic: *interaction with medium*

$$m \sim rL^3 \qquad \mathbf{V} \sim \mathbf{h}L$$

# Swimming and Length Scale

- Large Scales

Inertial  $t \rightarrow -t \Rightarrow F \rightarrow F$

? **constructive**

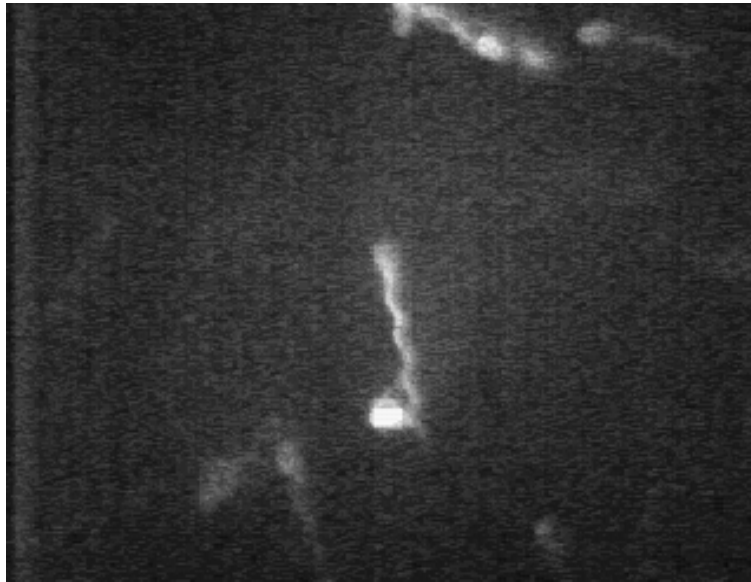
- Small Scales

Viscous  $t \rightarrow -t \Rightarrow F \rightarrow -F$

? **destructive**

# Bacterial Swimming

- *E. coli* driven by rotating flagella



L. Turner, W.S. Ryu, and H.C. Berg,  
*J. Bacteriol.* 182, 2793-2801 (2000)

- G.I. Taylor (1951)

Swimming at low Reynolds number

# Swimming at Low Reynolds

- Purcell Swimmer (1976)

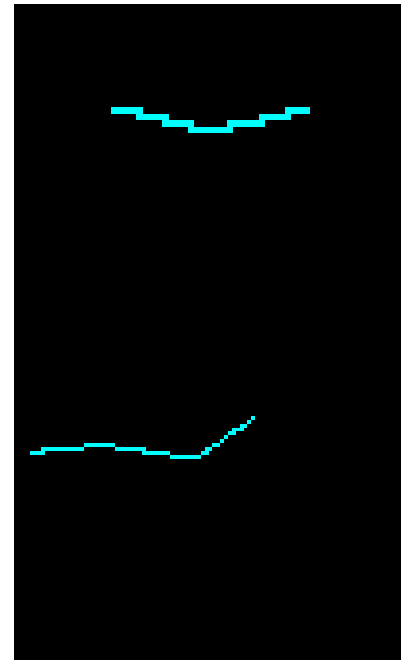
Using the smallest number of degrees of freedom:

- **One**, is not enough

*Scallop Theorem*

- **Two**, will just do!

*... but the passenger will get sea-sick!*



Courtesy of Daniel B. Murray

# Three-Sphere Swimmer

- Two Translational Degrees of Freedom

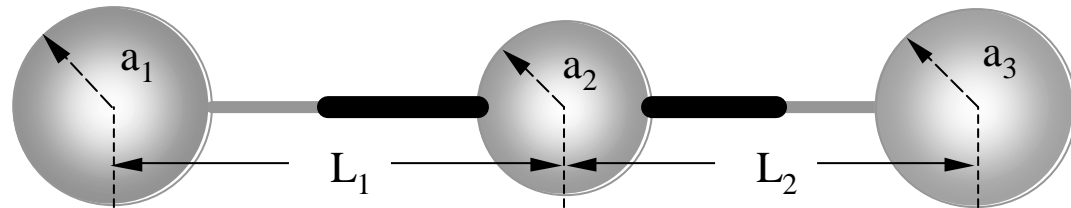


A. Najafi & R. Golestanian, PRE **69**, 062901 (2004)

# Swimming Velocity

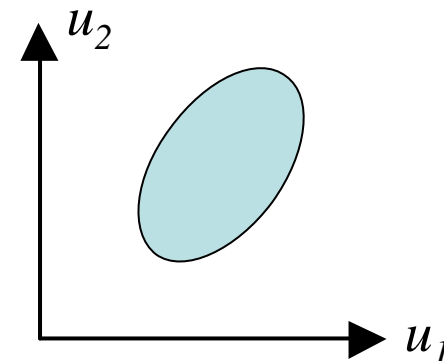
- Simple Perturbative Analysis

$$L_1 = \ell + u_1$$
$$L_2 = \ell + u_2$$



$$\bar{V} = \frac{7}{24} \frac{a}{\ell^2} \overline{(u_1 \dot{u}_2 - \dot{u}_1 u_2)}$$

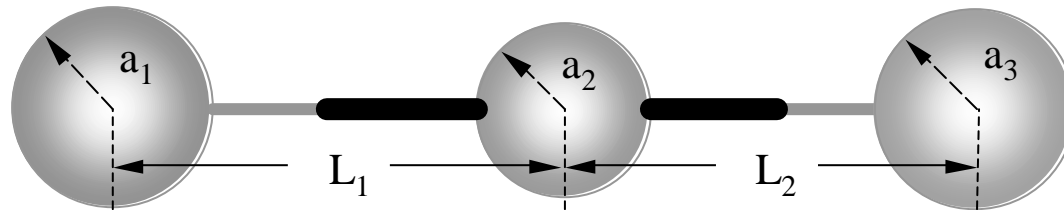
- Geometric Interpretation



# Harmonic Deformations

$$u_1(t) = d_1 \cos(\omega t + \varphi_1)$$

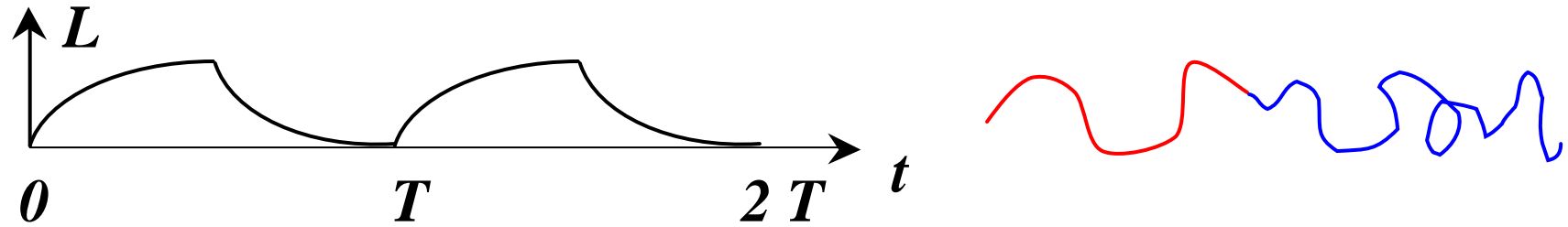
$$u_2(t) = d_2 \cos(\omega t + \varphi_2)$$



$$\overline{V} = \frac{7}{24} \frac{a}{\ell^2} d_1 d_2 \omega \sin(\varphi_1 - \varphi_2)$$



# Exponential Relaxation



$$\bar{V} = \frac{7}{96} \frac{a}{\ell^2} d_1 d_2 \left( \frac{1}{\tau_1} - \frac{1}{\tau_2} \right) \mathcal{F} \left( \frac{T}{4\tau_1}, \frac{T}{4\tau_2} \right)$$

$$\mathcal{F}(x, y) = \frac{1}{\sinh x \sinh y} \left[ \frac{\sinh(x+y)}{(x+y)} - \frac{\sinh(x-y)}{(x-y)} \right]$$

# Noisy Deformations

$$\overline{V} = \frac{7}{24} \frac{a}{\ell^2} \frac{1}{T} \int \frac{d\omega}{2\pi} i\omega \overline{[u_2(\omega)u_1(-\omega) - u_1(\omega)u_2(-\omega)]}$$

$$u_i(t) = \sum_n d_{in} \cos(\omega_n t + \varphi_{in})$$

$$\overline{V} = \frac{7}{24} \frac{a}{\ell^2} \sum_n \overline{d_{1n}d_{2n}\omega_n \sin(\varphi_{1n} - \varphi_{2n})}$$

# Load-Velocity Response

- Force-Velocity Relation

$$\bar{V}(F) = \frac{7}{24} \frac{a}{\ell^2} \overline{(u_1 \dot{u}_2 - \dot{u}_1 u_2)} + \frac{F}{18\pi\eta a}$$

- Stall Force

$$F_s = -18\pi\eta a \bar{V}(0)$$

# Mechanical Efficiency

- Power Consumption

$$\overline{\mathcal{P}} = 2\pi\eta a \left[ \overline{\dot{u}_1^2} + \overline{\dot{u}_2^2} + \overline{(\dot{u}_1 + \dot{u}_2)^2} \right]$$

- Lighthill's Efficiency  $\eta_L = \frac{18\pi\eta a \overline{V}^2}{\overline{\mathcal{P}}}$

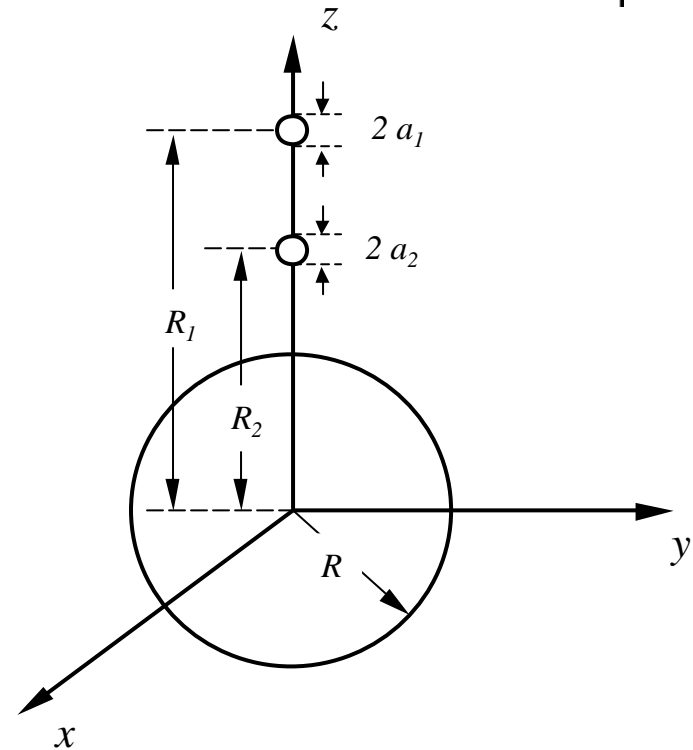
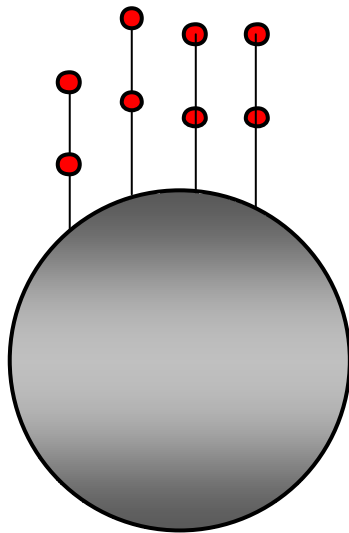
$$\eta_L = \frac{49 a^2}{64 \ell^4} \frac{\overline{(u_1 \dot{u}_2 - \dot{u}_1 u_2)^2}}{\overline{\dot{u}_1^2} + \overline{\dot{u}_2^2} + \overline{(\dot{u}_1 + \dot{u}_2)^2}}$$

# Swimmer with Cargo

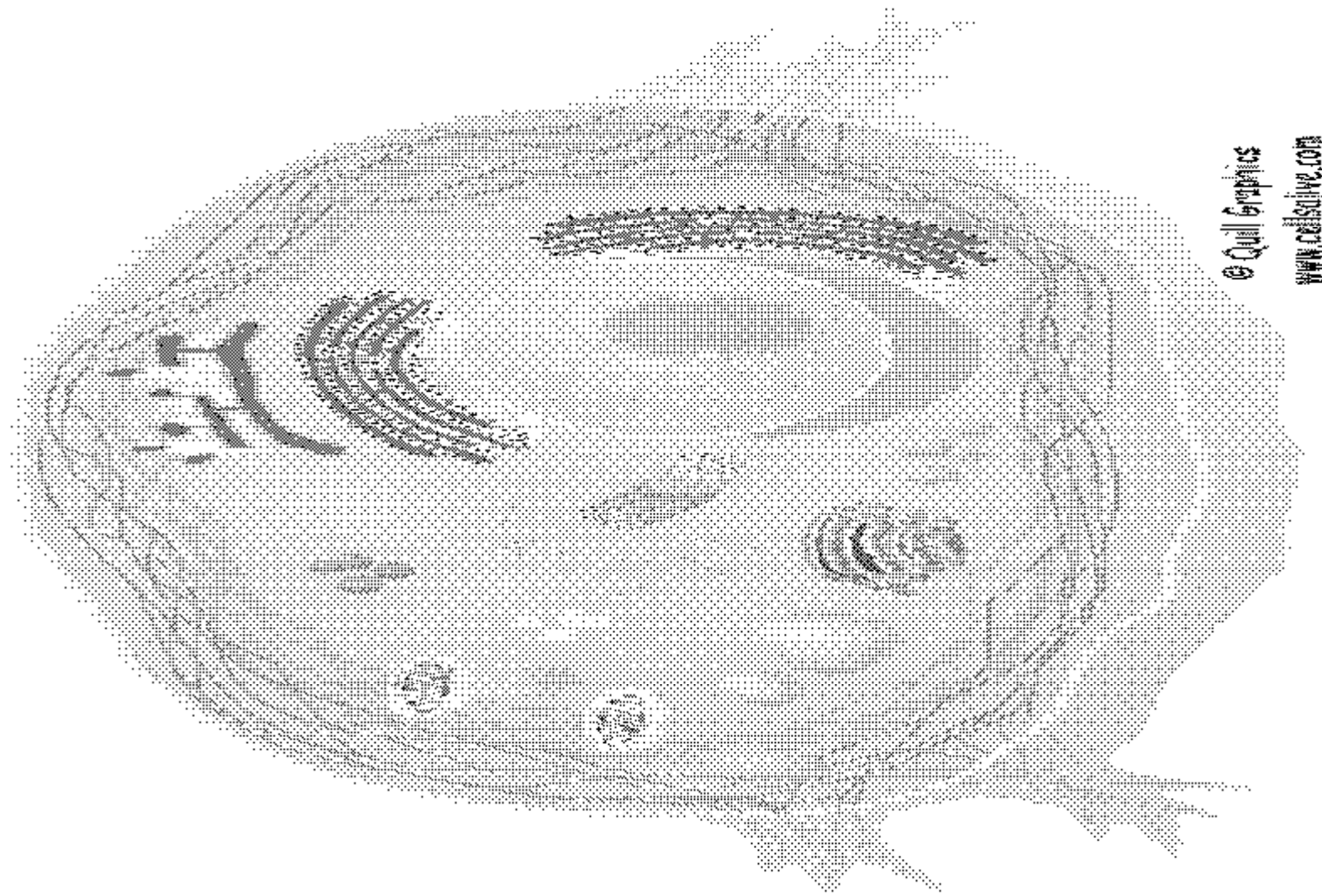
- Small deformations

$$\overline{V} = \frac{9}{4} \left( \frac{a_1 a_2}{R^3} \right) \frac{\ell_2^2 (3\ell_1 + 2\ell_2)}{\ell_1^2 (\ell_1 + 2\ell_2)} \frac{(\dot{u}_1 u_2 - \dot{u}_2 u_1)}{\left| \right.}$$

- Optimal Design



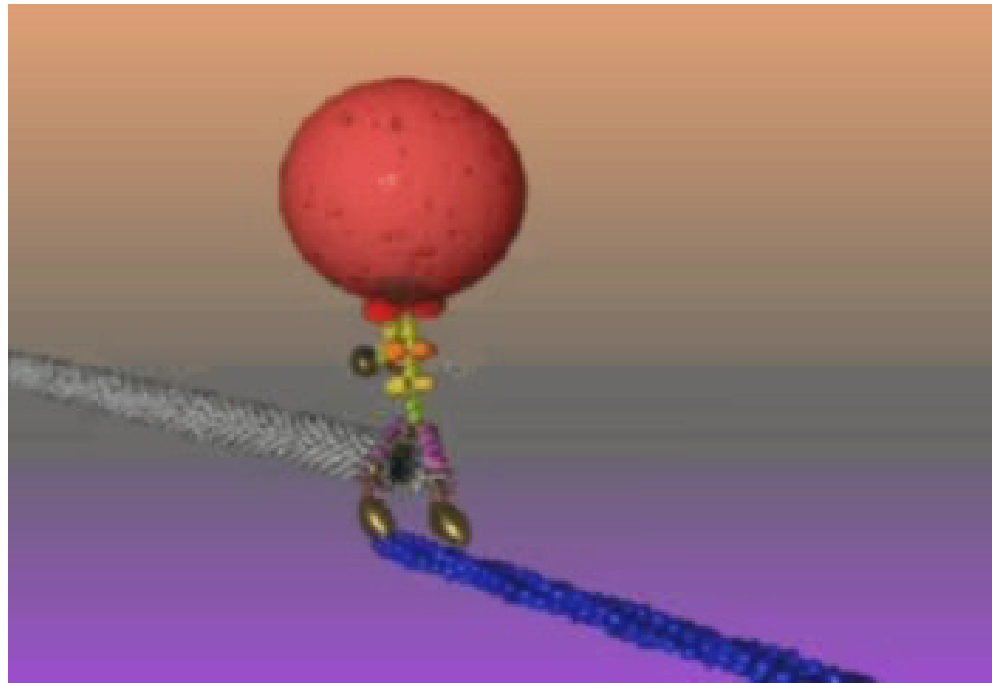
# Transport System in the Cell



© Quill Graphics  
www.cellshive.com

# Molecular Motors

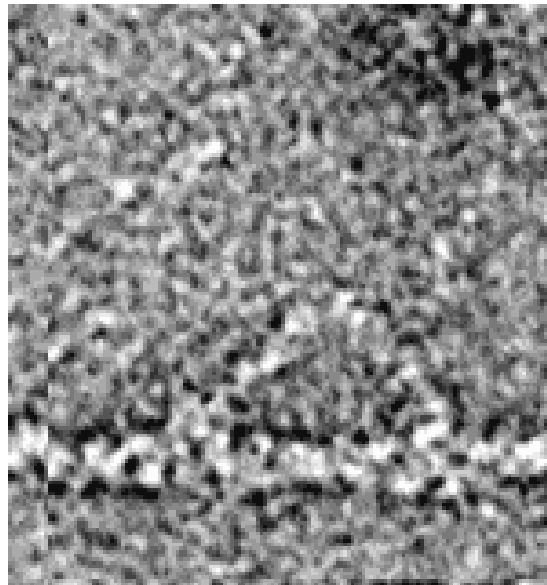
- Motion of myosin-V molecules on actin



Courtesy of Dr. George Langford

# Molecular Motors

- Myosin-V: ``Electron Video Microscopy''



The Muscle Group, Leeds 2000

M L. Walker, S A. Burgess, J R. Sellers,  
F Wang, J.A. Hammer, J Trinick, & P J.  
Knight, Nature **405**, 804-807 (2000)



# Molecular Motors

- Principal Characteristics
  - Nano-scale machines
  - Highly accurate despite Brownian agitation
  - Converting *chemical* energy directly to useful *mechanical* work

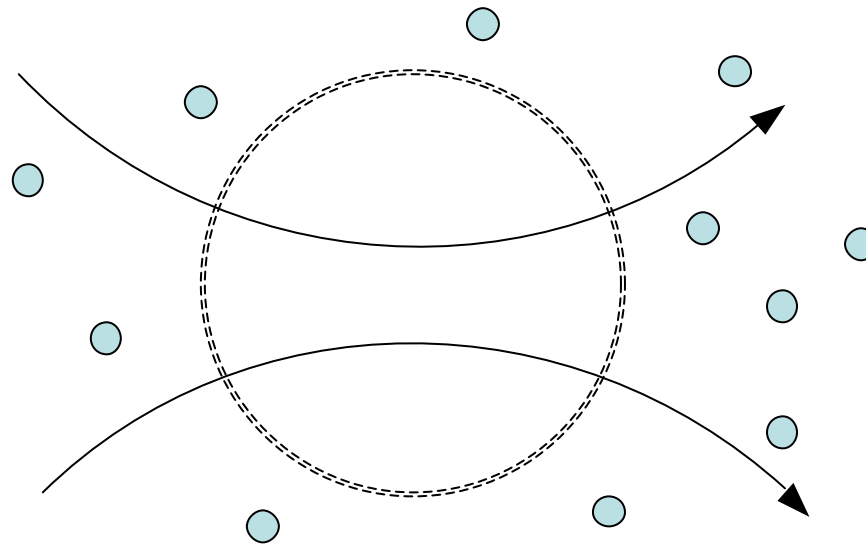
Can we design such a machine?

# Phoretic Locomotion

- Osmio-phoresis

Motion of lipid vesicles under a gradient of sugar molecules

? Nardi, Sackmann & Bruinsma, PRL **82**, 5168 (1999)

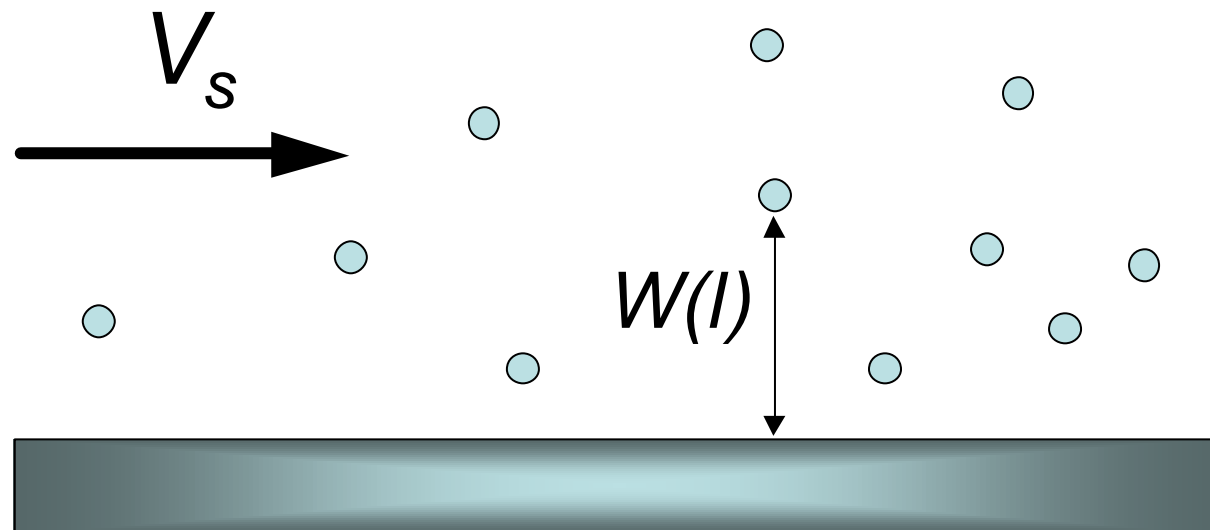


# Phoretic Locomotion

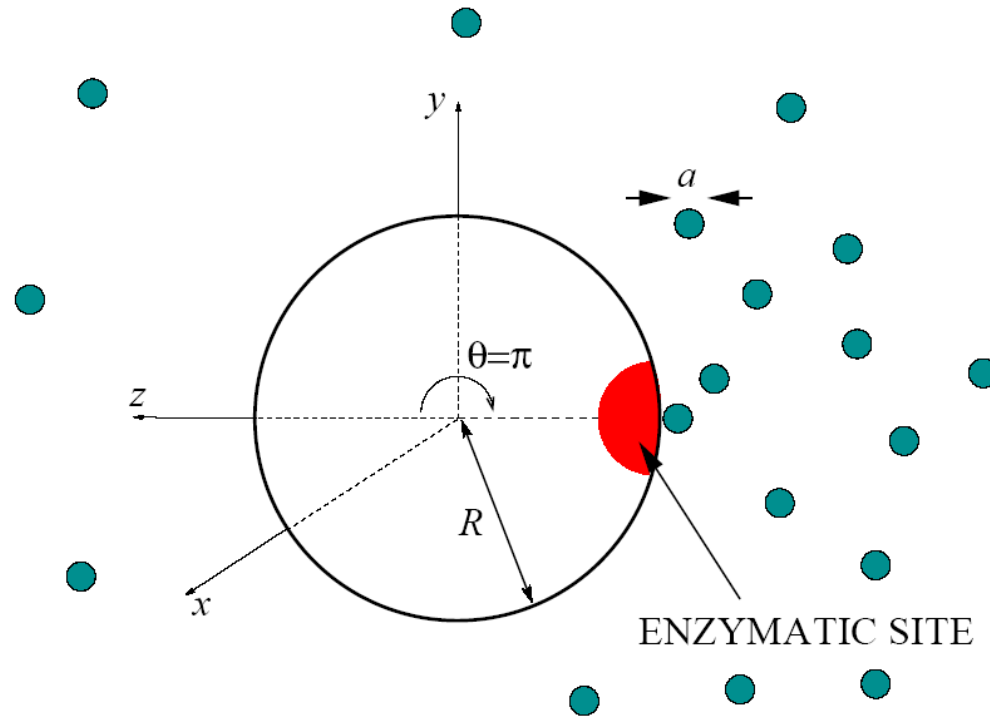
- Diffusio-phoresis

? Derjaguin, Dukhin, Korotkova, Kolloidn. Zh. **23**, 53 (1961)

? Anderson & Prieve, Sep. Purif. Methods **13**, 67 (1984)



# Reaction-Driven Propulsion



R. Golestanian, T.B. Liverpool & A. Ajdari, PRL **94**, 220801 (2005)

# Propulsion Velocity

$$v_0 = \frac{\ell}{\tau_f} \quad \ell = a \left( \frac{3 \lambda^2}{4 R^2} \right)$$

- Diffusio-phoresis

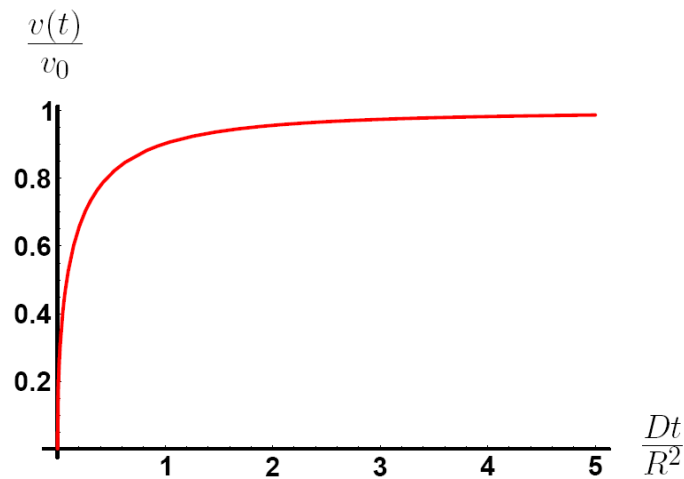
$$\lambda_D^2 = \int_0^\infty dl \, l \left[ 1 - e^{-W(l)/k_B T} \right]$$

- Osmio-phoresis

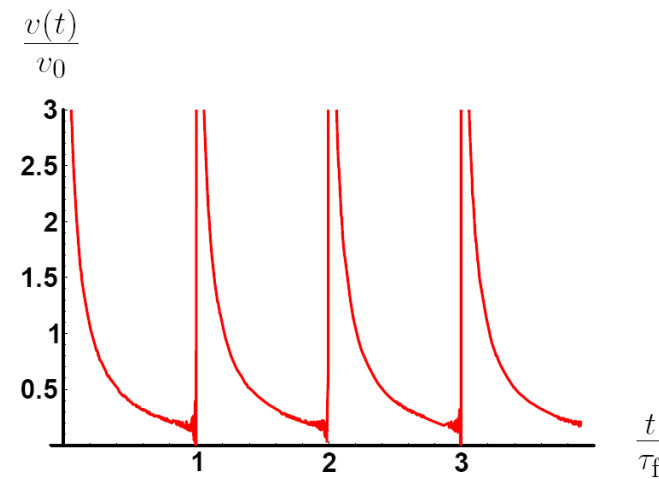
$$\lambda_O^2 = R^2 \left[ \frac{\eta L_p / R}{2 + 20 \eta L_p / R} \right]$$

# Time Response

- Relaxation:



- Periodic particle release:

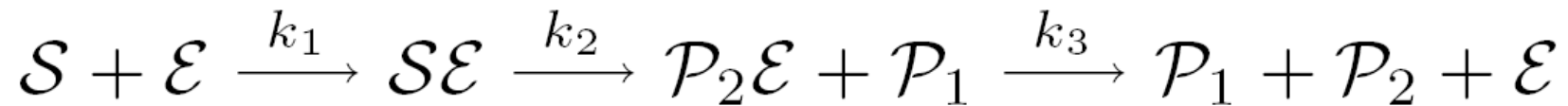


- Fluctuations:

$$\frac{(\Delta v)^2}{v_0^2} \sim D\tau_f / R^2$$

# Reaction Kinetics

- Michaelis-Menten Rule



$$\frac{1}{\tau_f} = \left( \frac{k_2 k_3}{k_2 + k_3} \right) \frac{C_S(\mathcal{E})}{K_M + C_S(\mathcal{E})}$$

Michaelis constant:  $K_M \equiv \frac{1}{k_1} \left( \frac{k_2 k_3}{k_2 + k_3} \right)$

# Brownian Agitation

- Rotational Diffusion Time

$$\tau_R = 8\pi\eta R^3 / (k_B T)$$

Directional persistence increases with size



# Optimal Design

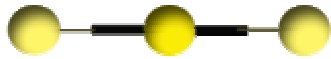
- Vesicle/bead of radius:  $R = 2 \mu\text{m}$
- Acetylcholinesterase:  $1/\tau_f \simeq 25000 \text{ s}^{-1}$
- Propulsion velocity:  $v_0 \sim 0.5 \mu\text{m/s}$
- Rotational diffusion time: 50 s

? *Directional Locomotion*

# Conclusion

- Model A:

? externally controlled



- Model B:

? reaction-driven

